

Driven complex matter

Paul van Loosdrecht
University of Cologne



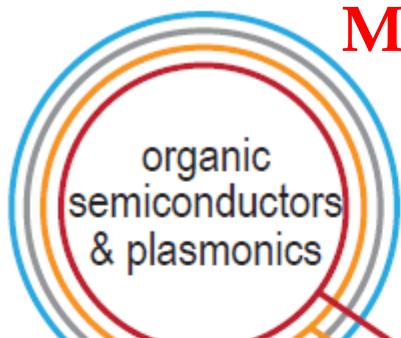
UoC: Excellence University

The University of Cologne's **institutional strategy “Meeting the Challenge of Change and Complexity”** aims at strengthening and further developing the university's research profile, at establishing funding programmes for cutting edge research and at integrating new career-promoting structures and support measures. The strategy also includes the further development of regional and international research networks and exchange programs, the promotion of Gender Equality and a set of measures to promote research-based teaching.

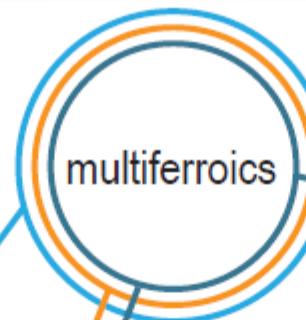
QM2 Quantum matter and materials

“Quantum Matter and Materials” (QM2) is a fascinating field of research driven both by the intellectual challenge and the promise for applications. Within QM2, researchers of mathematics, experimental and theoretical physics, anorganic and physical chemistry and crystallography collaborate to unravel the properties of quantum matter. The mathematical structures underlying topological matter, the prospect to functionalize “Dirac matter” like graphene, new states of matter arising from spin-orbit interactions, quantum matter far from thermal equilibrium, the use of nanostructured materials for wide ranges of applications, and the development of organic electronics, are some of the research topics important for QM2.

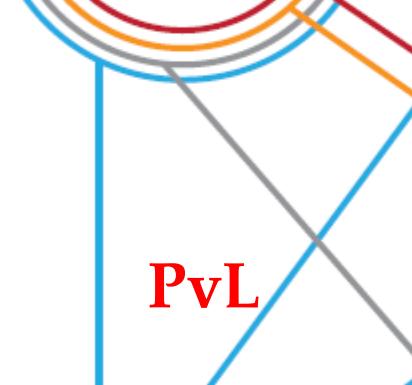
chemistry



Meerholz



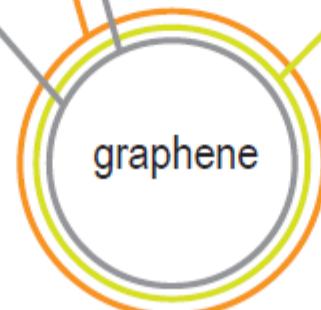
Braden
Grueninger
PvL



PvL

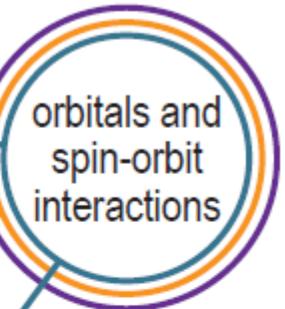


Michely

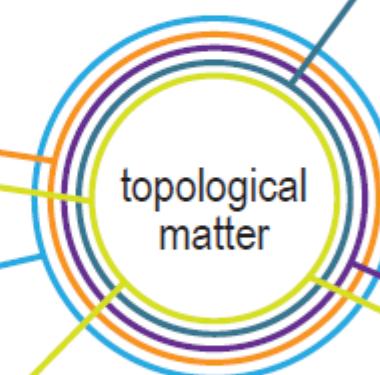


mesoscopics

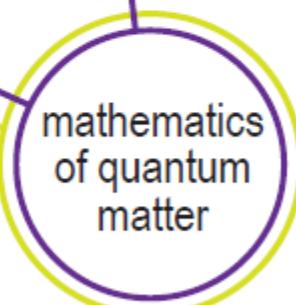
correlated
matter



Trebst



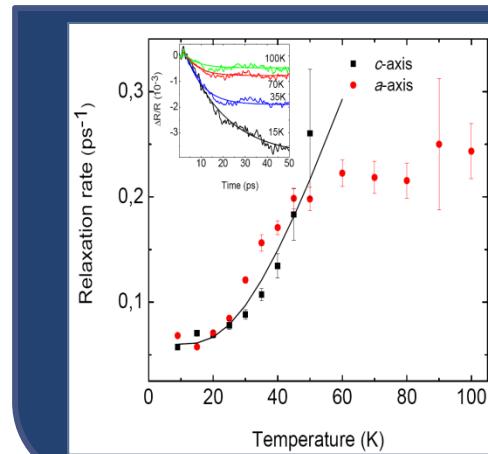
Altland



Dynamical properties of matter

Mapping interactions & processes on the time domain

- Electron-electron
- Electron-phonon
- Electron-magnon
- Phonon-magnon
- Energy, charge, and spin transfer
- Spin, phonon, orbital coherence
- Pair breaking
- Charge separation
- ...



Magnon assisted hopping in TbMnO_3

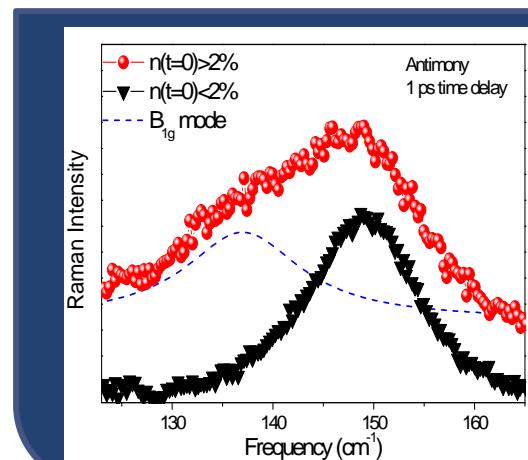
- Charge – Magnon interaction
- $k=\pi$ Electromagnon energy

I.P. Handayani et al., JPCM 2013

Dynamical properties of matter

Controlling states of matter

- Novel states of matter
- Ultrafast melting of lattice,
spin, charge ordering
- Switching magnetism
- Switching ferroelectricity
- Stripes & induced
superconductivity
- Metastable states
- ...



Induced phase transition in A3 elemental metals

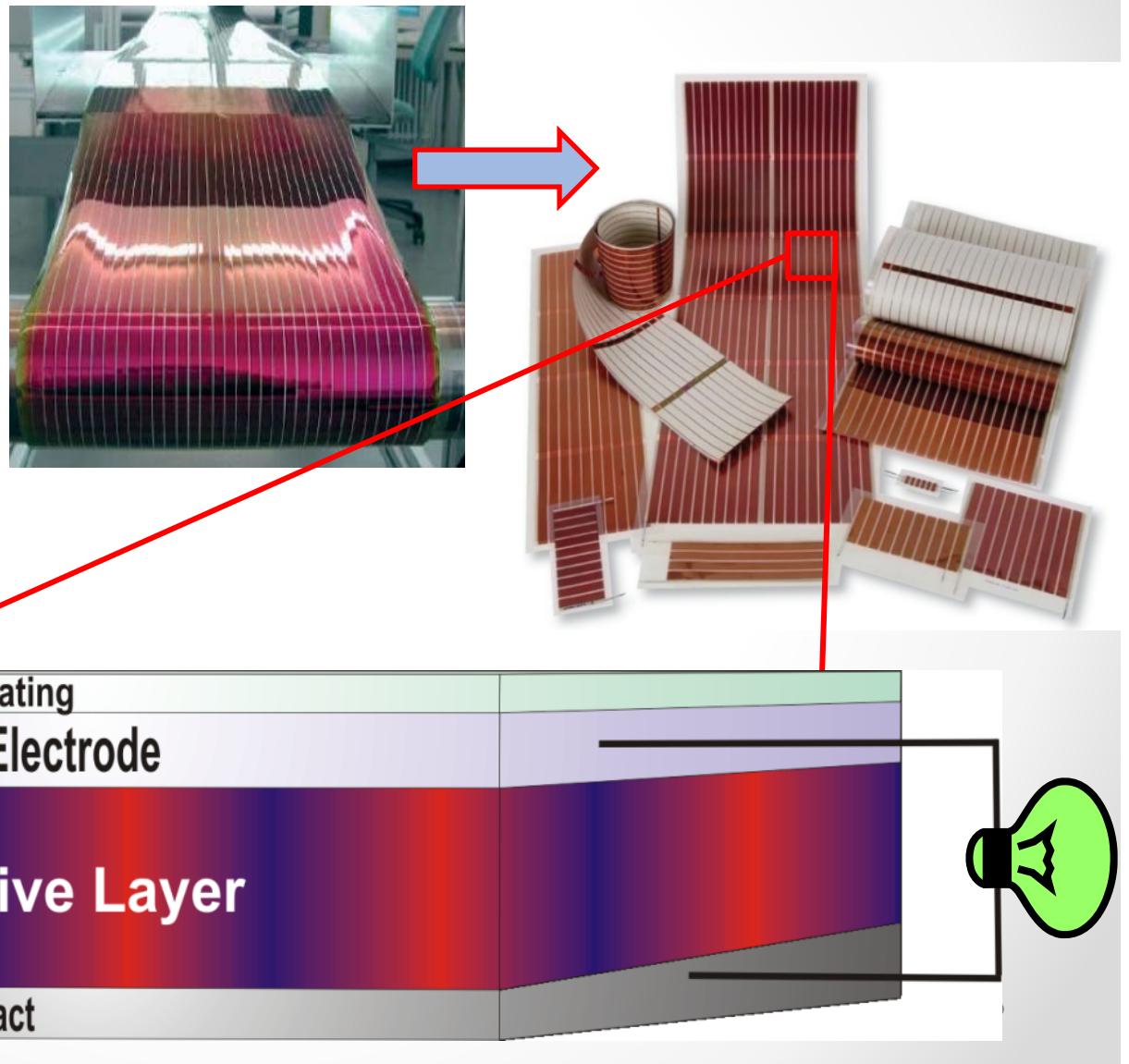
- Non-thermal state
- Destabilization Jones-Peierls state
- New Jones-Peierls distortion

D. Fausti et al., PRB-RC 2009

Charge separation in OPV

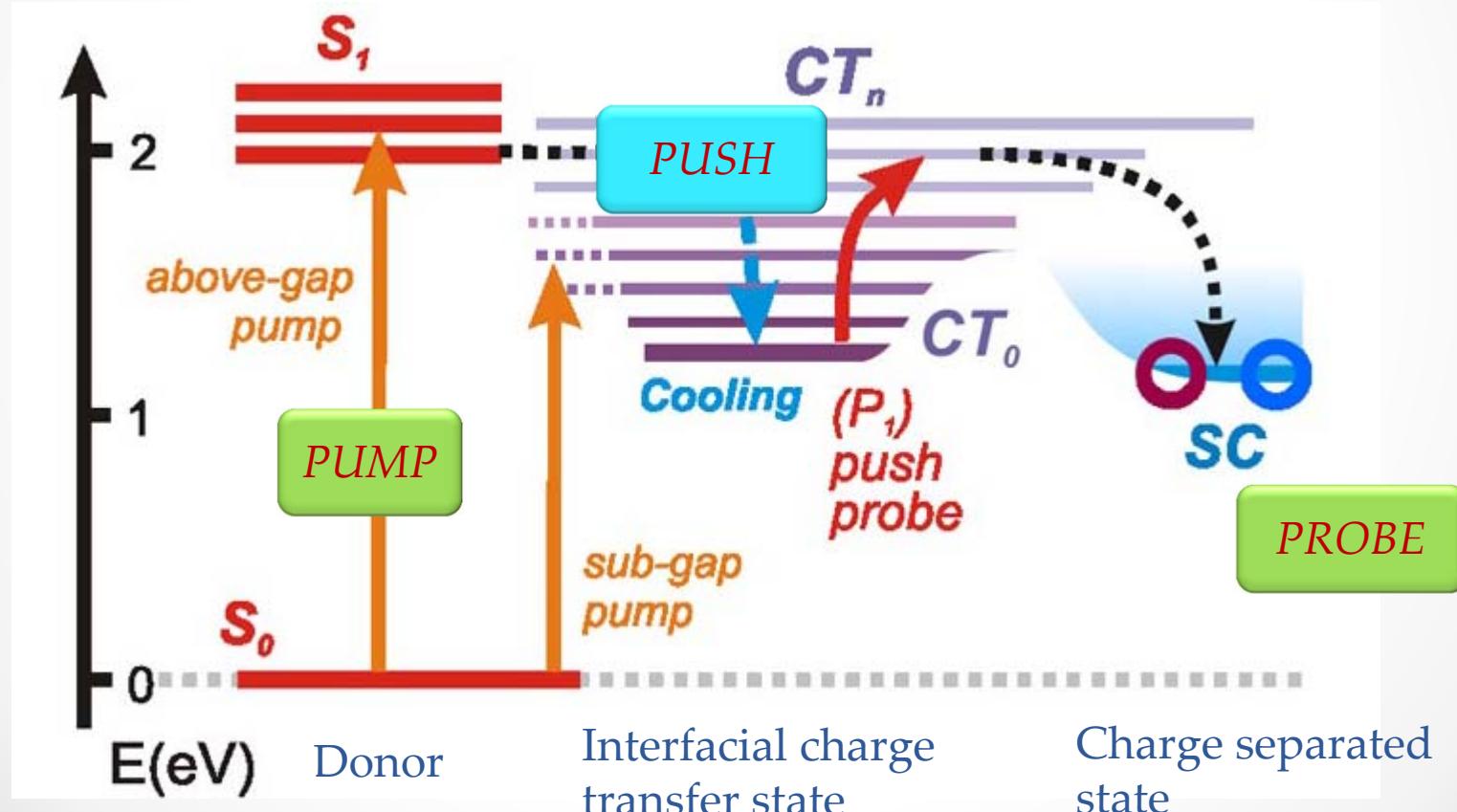
Advantages:

- ✓ Inexpensive
- ✓ Lightweight
- ✓ Flexible
- ✓ Tunable properties
- ✓ Printing-processable



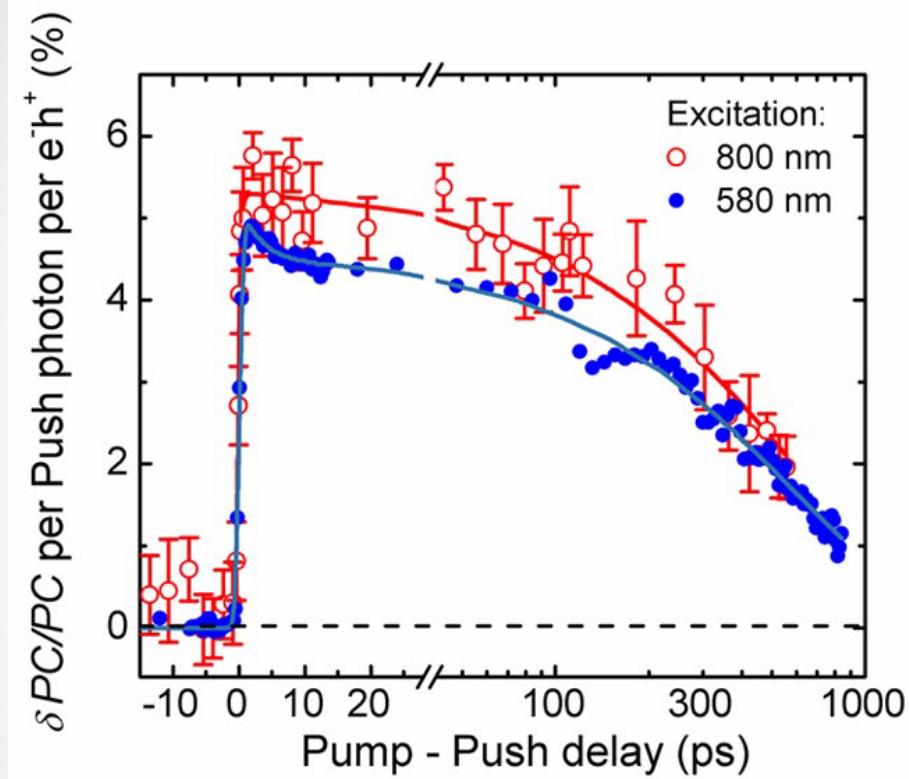
Charge separation in OPV

*Action occurs at the donor-acceptor interface.
What determines the efficiency ?*

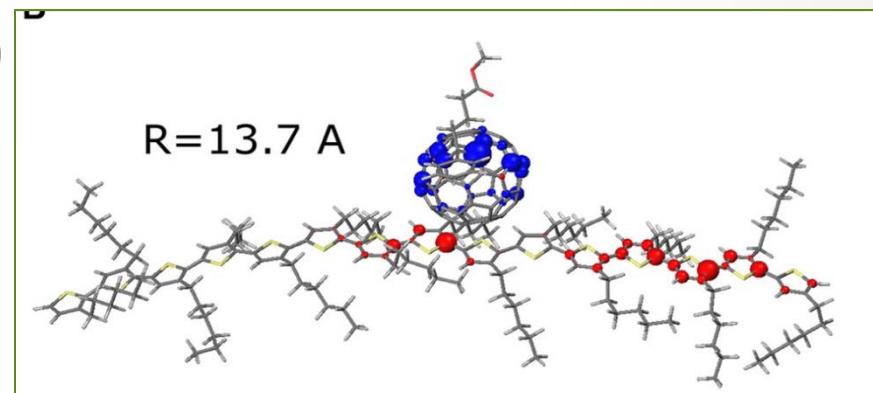


Similar for acceptor: Bakulin et al., Funct. Adv. Mat. 2010

Charge separation in OPV



*Delocalization drives
charge generation efficiency*



Ultrafast magnetism in EuO

Interplay between magnetism and conductivity

RKKY interaction

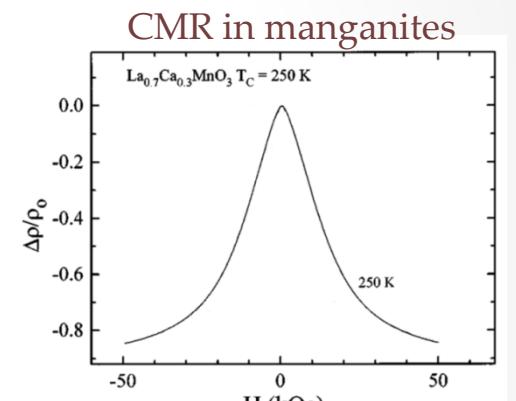
- GMR read heads
- Kondo physics

Double exchange

- CMR in manganites
- Heavy fermion systems

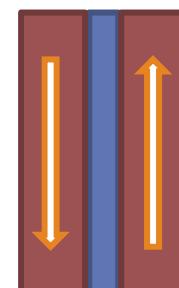
Spintronics

- Doped semiconductors
- EuO



M.F. Hundley *et al.* APL 67, 860 (1995)

spin valve



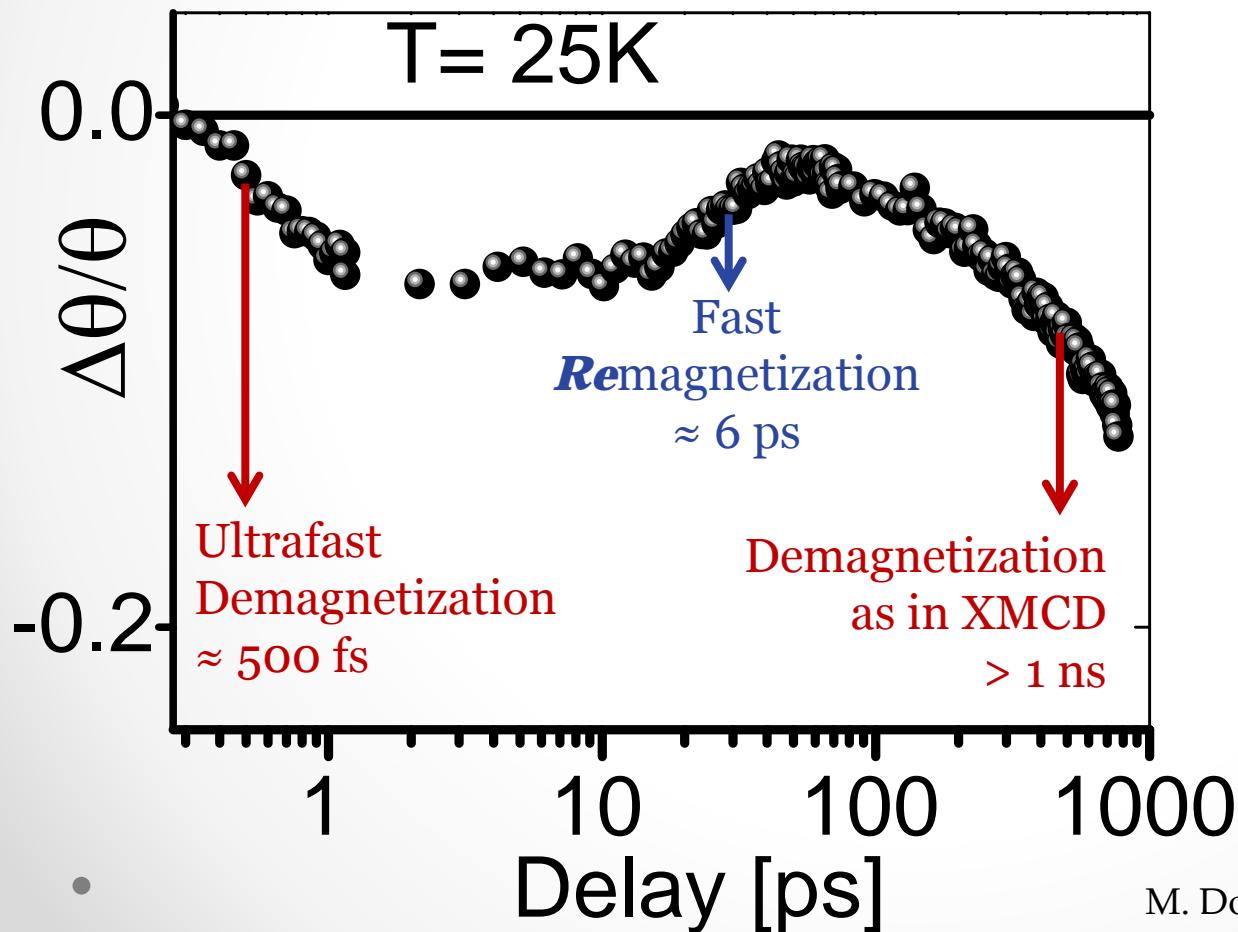
GMR

hard drive



Ultrafast magnetism in EuO

Ultrafast switching of magnetism in a magnetic semiconductor:
Interplay between 5d¹ electrons, 4f⁶ holes and 4f⁷ moments



First picosecond
hot 5d¹ electrons
Elliot-Yafet scattering

10 ps timescale
4f⁶-5d¹
magnetic polaron
formation

Nanosecond timescale
4f⁶,4f⁷
spin-lattice relaxation
anisotropic CF
fluctuations

Dynamic heat transport in low dimensional quantum magnets

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University of Cologne
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Universität zu Köln





university of
groningen

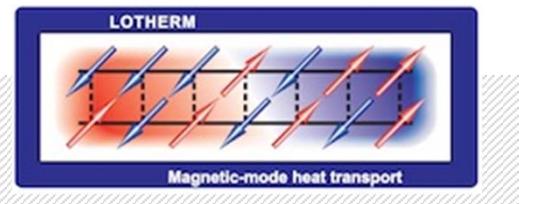


Dynamic heat transport in low dimensional quantum magnets.

Paul H.M. van Loosdrecht

Optical Condensed Matter Physics

Zernike Institute for Advanced Materials



The LOTHERM collaboration

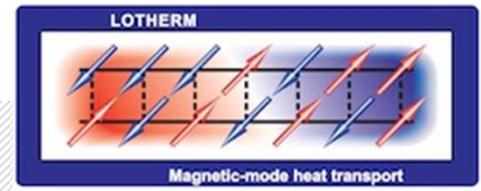


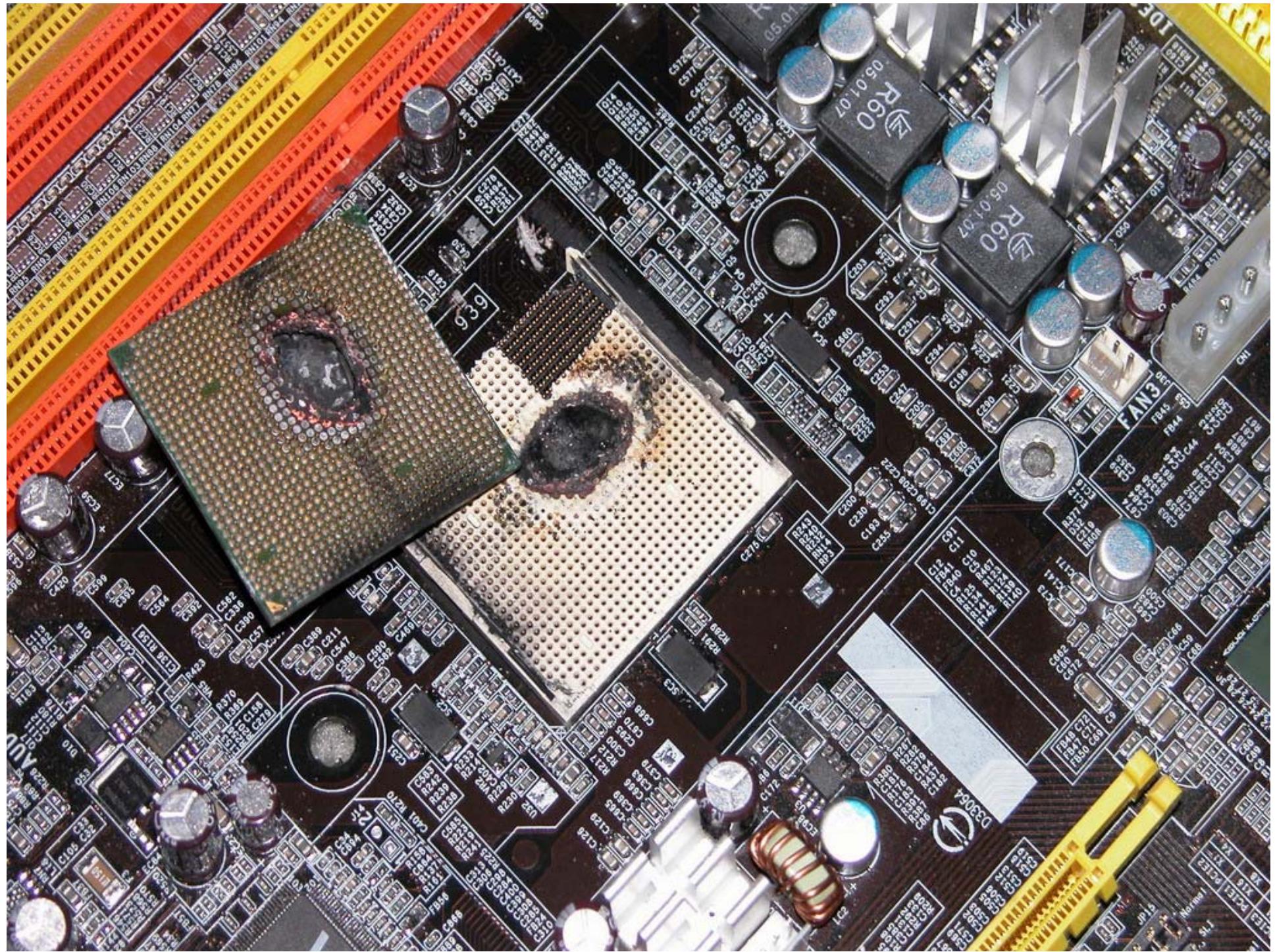


Heat management

“Managing heat generation within IC’s will be a crucial issue in developing the next generation of electronics.”

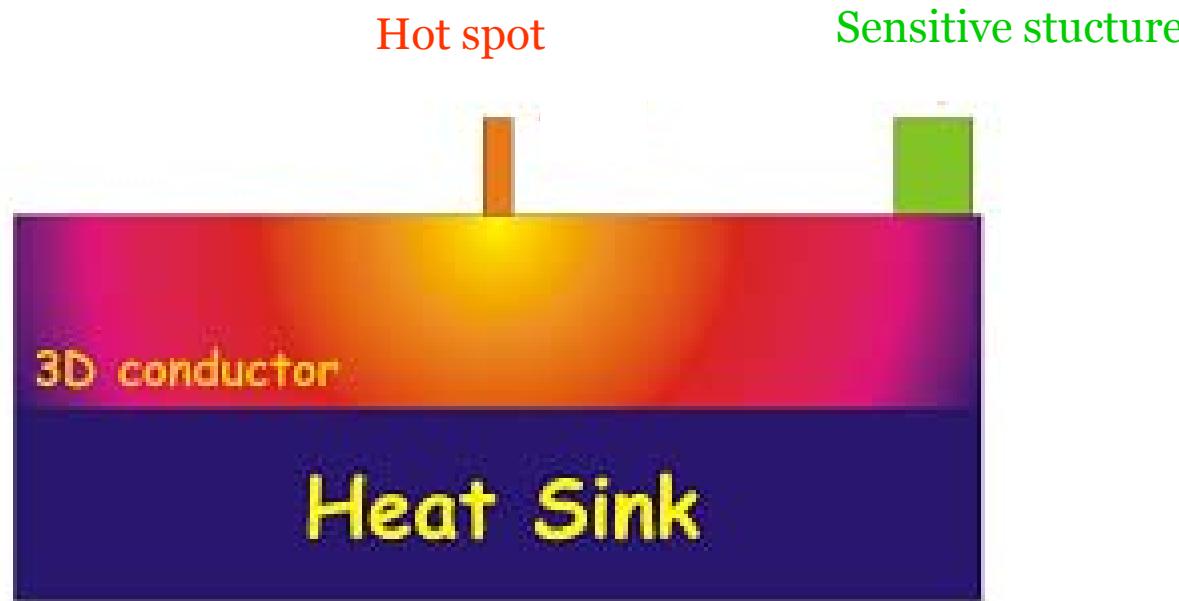
International Technology Roadmap for Semiconductors, ITRS-2009







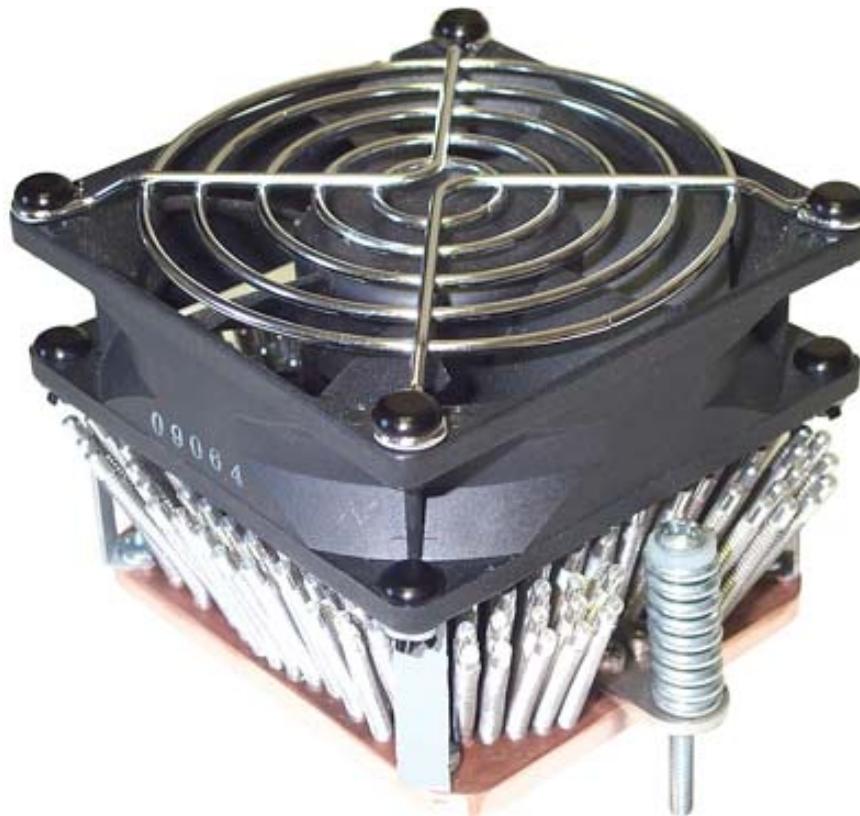
Heat management



Solutions ?



Buy a big fan





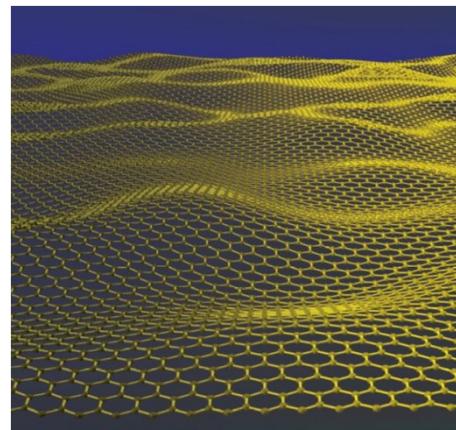
Use a champion heat conductor

Diamond
2000 W/mK



Hope

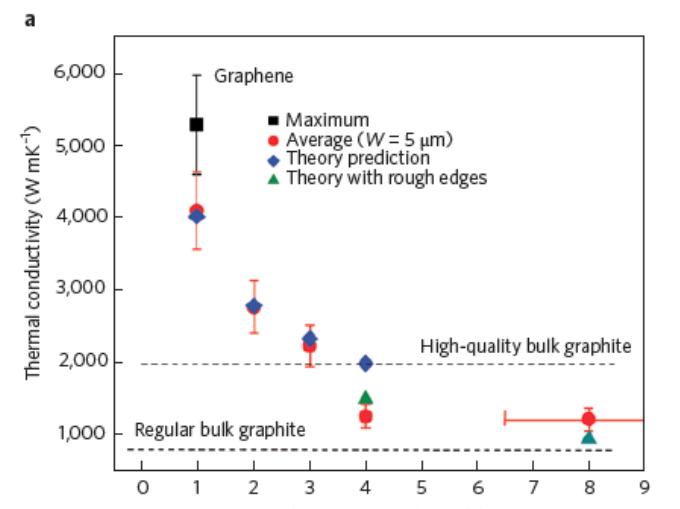
Graphene 5000 W/mK



Even more Hope

Expensive

Also good charge conductor

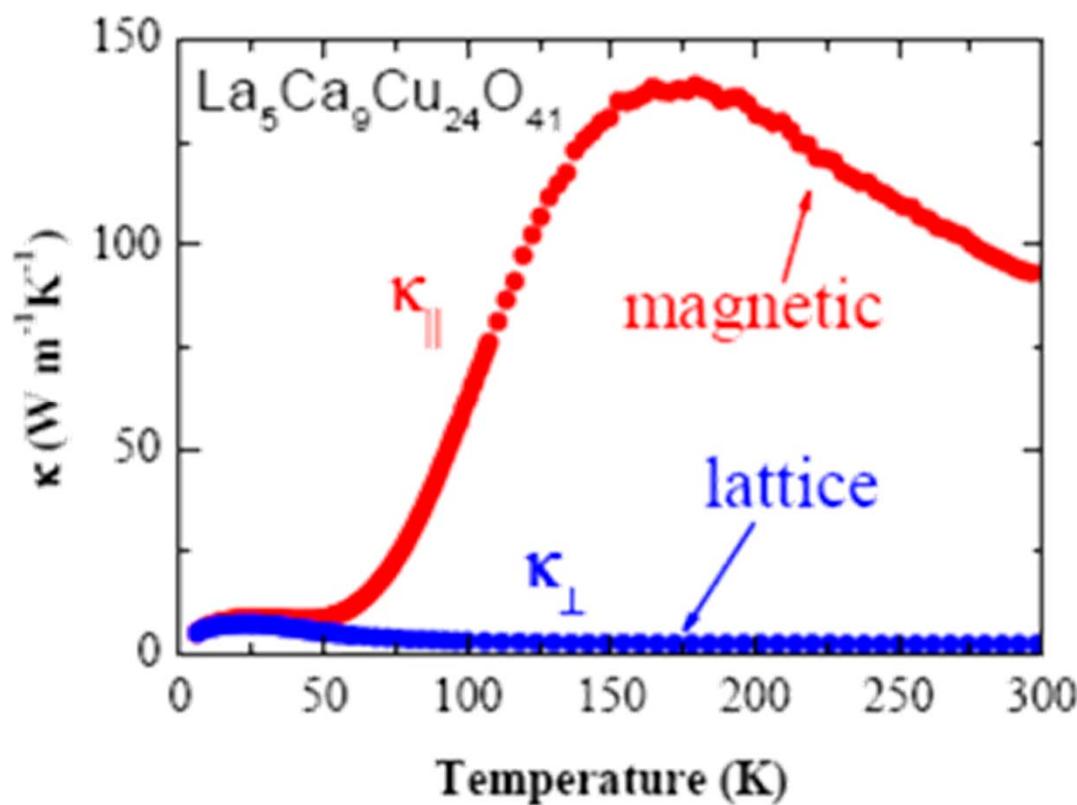




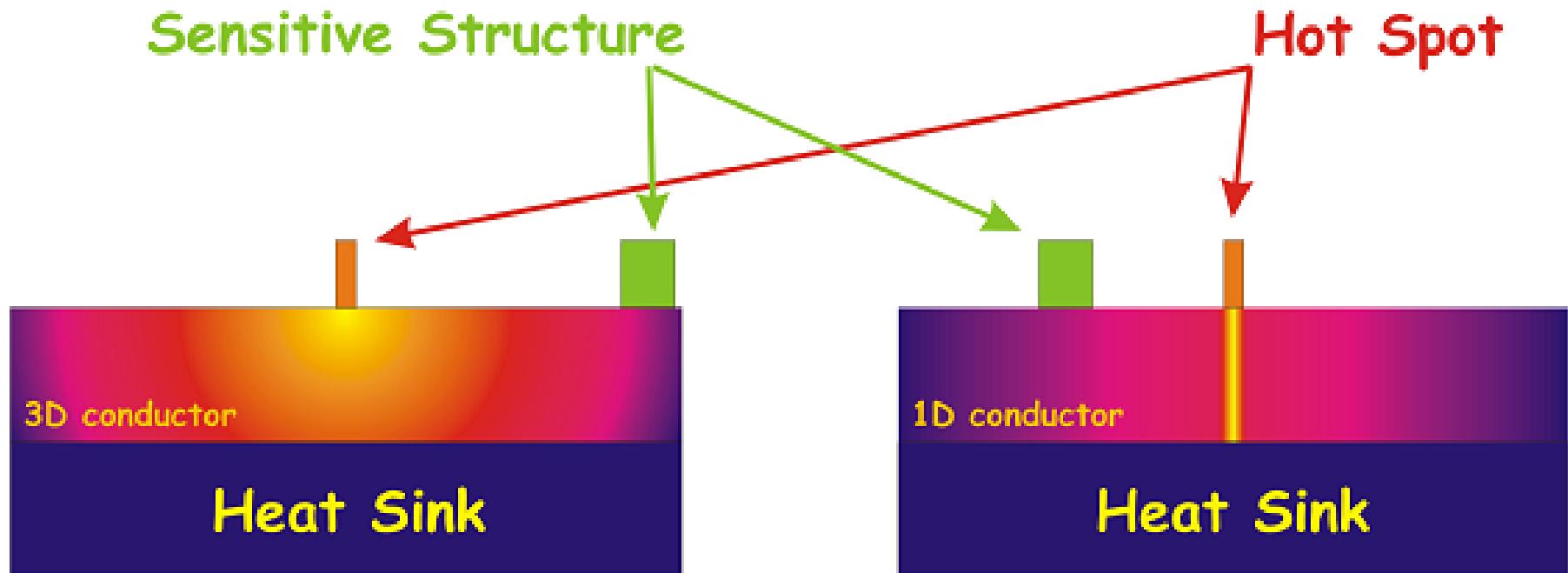
Quantum magnets ?

LCCO:

- 140 W/mK
- Good charge insulator
- 1 dimensional



Steel	50 W/mK
Copper	350 W/mK
Silicon	100 W/mK
SiO_2	1 W/mK
SiC	400 W/mK
Diamond	1000 W/mK

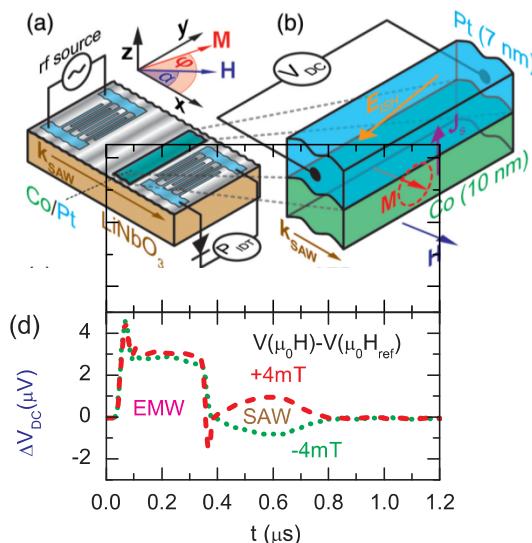


Low dimensional quantum magnets

- Good heat conductivity ($> 100 \text{ Wm}^{-1}\text{K}^{-1}$)
- Unidirectional
- Electrically insulating



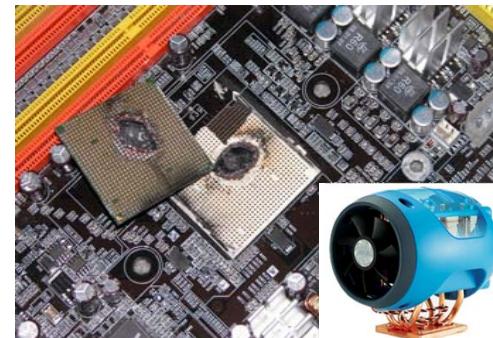
Spin pumping with coherent elastic waves



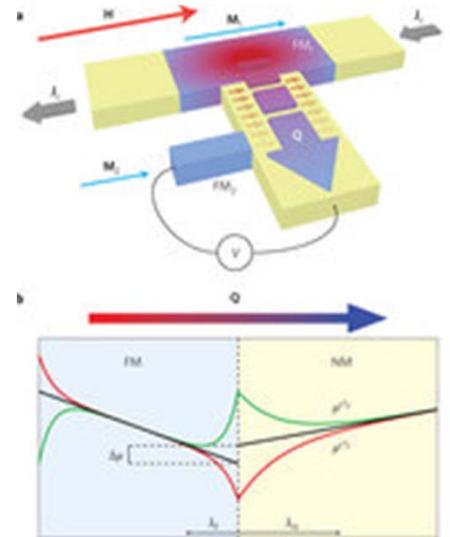
Weiler, PRL 2012

PHONON-MAGNON COUPLING

Heat management



Spin Seebeck Effect



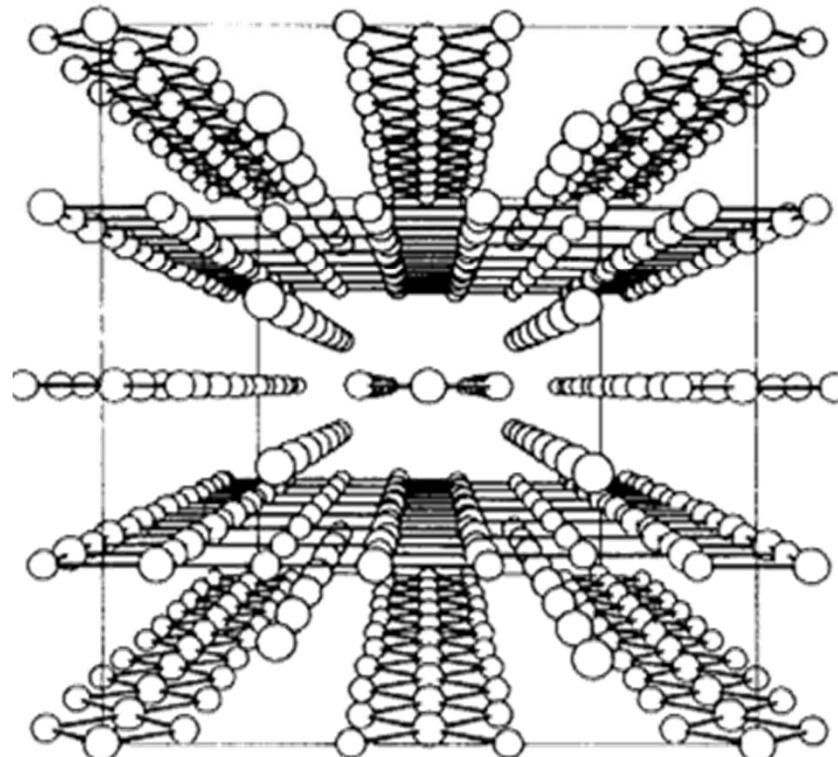
Bauer, Nat. Mat. 2012

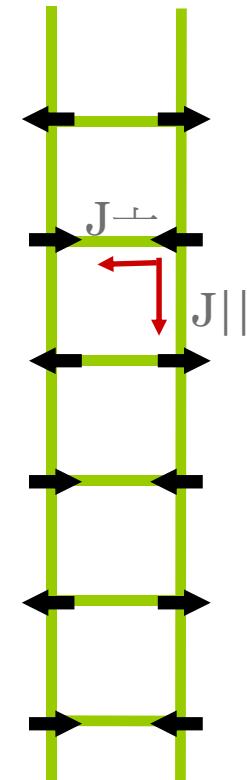
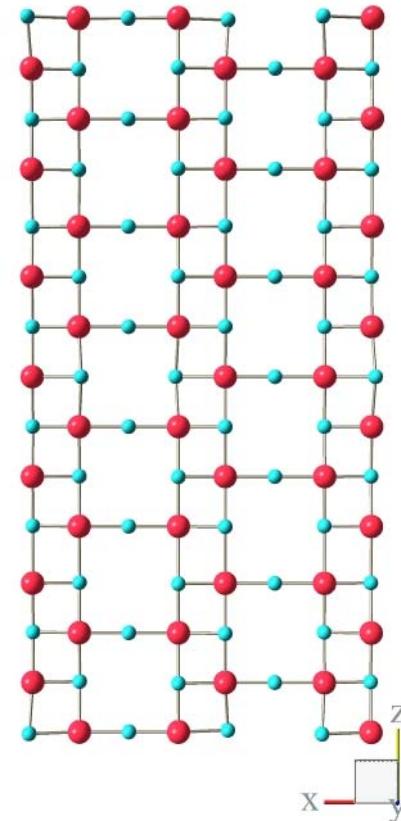
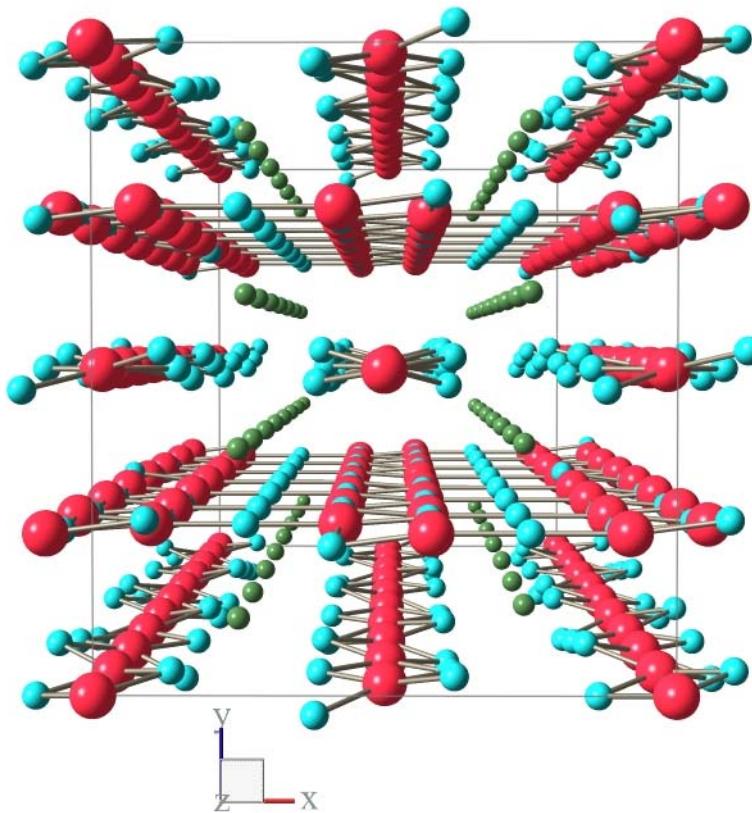


Mat. Res. Bull., Vol. 23, pp. 1355-1365, 1988. Printed in the USA.
0025-5408/88 \$3.00 + .00 Copyright (c) 1988 Pergamon Press plc.

THE INCOMMENSURATE STRUCTURE OF (Sr_{14-x}Ca_x)Cu₂₄O₄₁ (0 < x ~ 8)
A SUPERCONDUCTOR BYPRODUCT

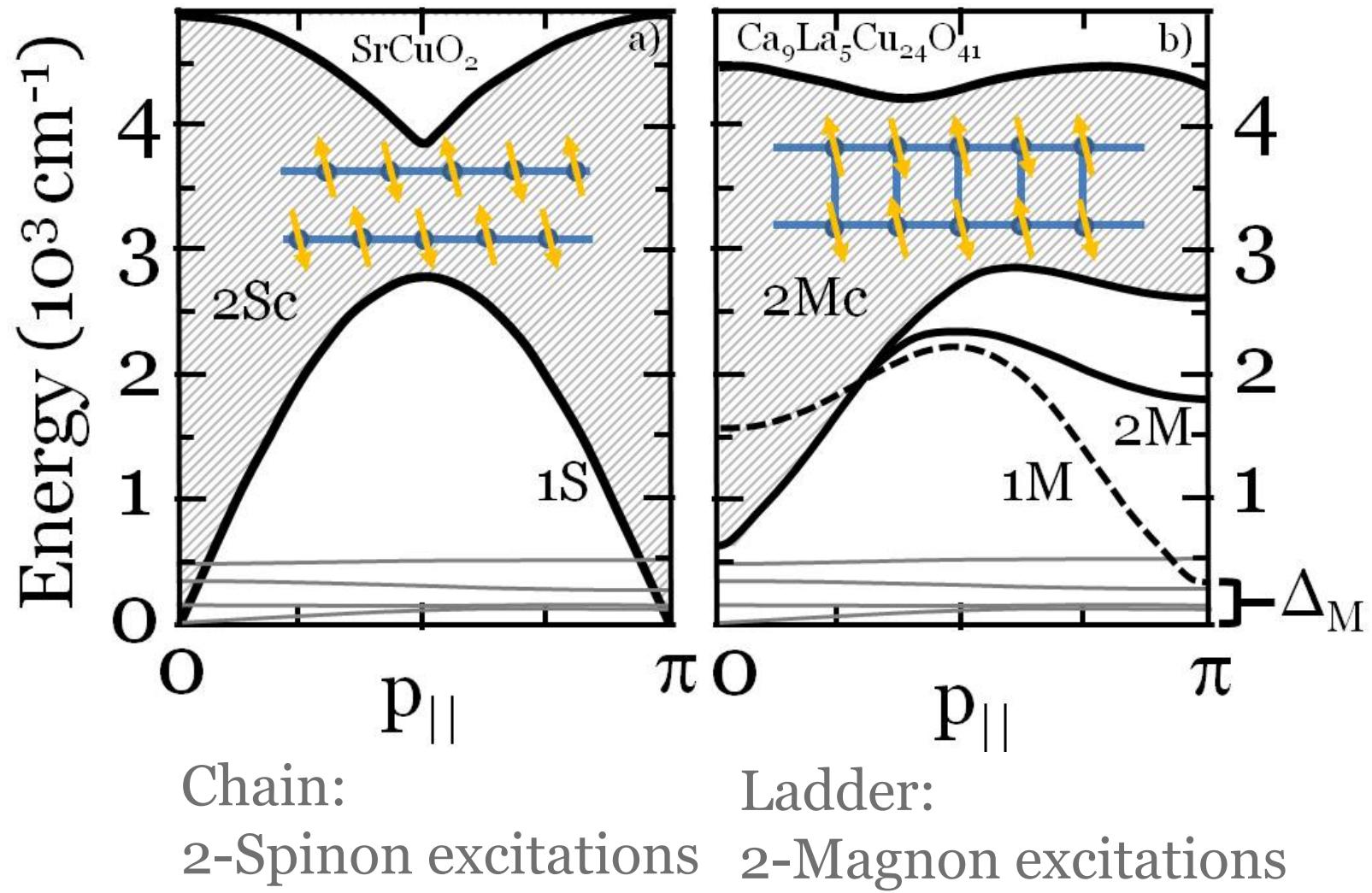
E. M. McCarron, III*, M. A. Subramanian, J. C. Calabrese
and R. L. Harlow*





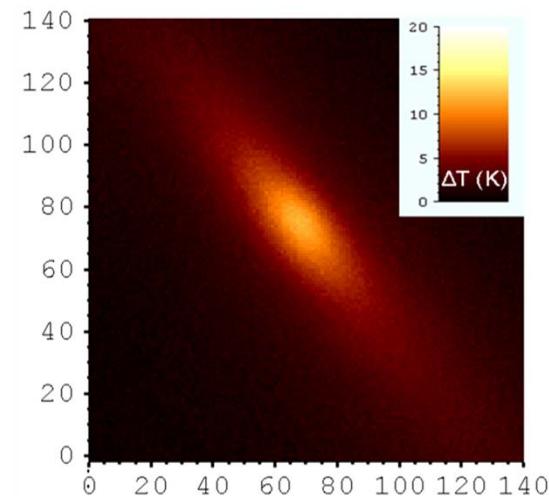
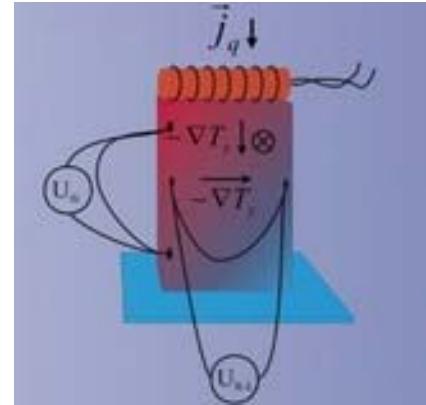
(a)

$$\begin{aligned}J_{//} &= 130 \text{ meV} \\J_{\perp} &= 70 \text{ meV} \\\Delta &= 32 \text{ meV}\end{aligned}$$





- › Traditional approach
 - Slow
 - Not suitable for films
- › 3ω approach
 - Not fast enough
 - Elaborate
- › Optical approach
 - Fast
 - Also on thin films
 - Easy



Otter et al., Jmmm 321, 796 (2009)

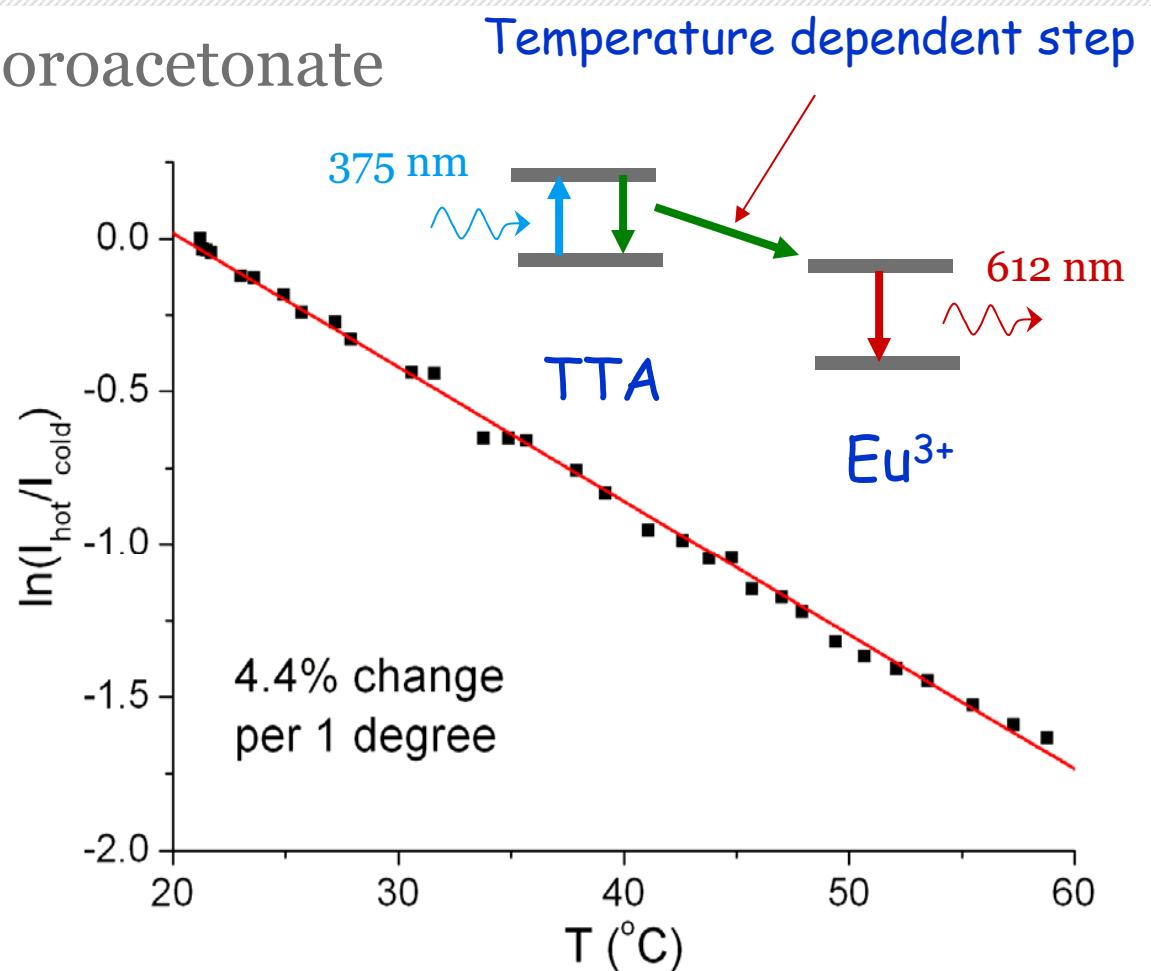
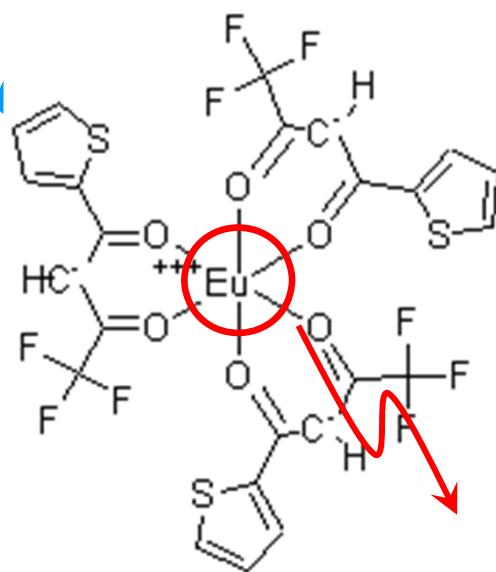


Measuring heat conductivity requires a heat source and temperature measurement

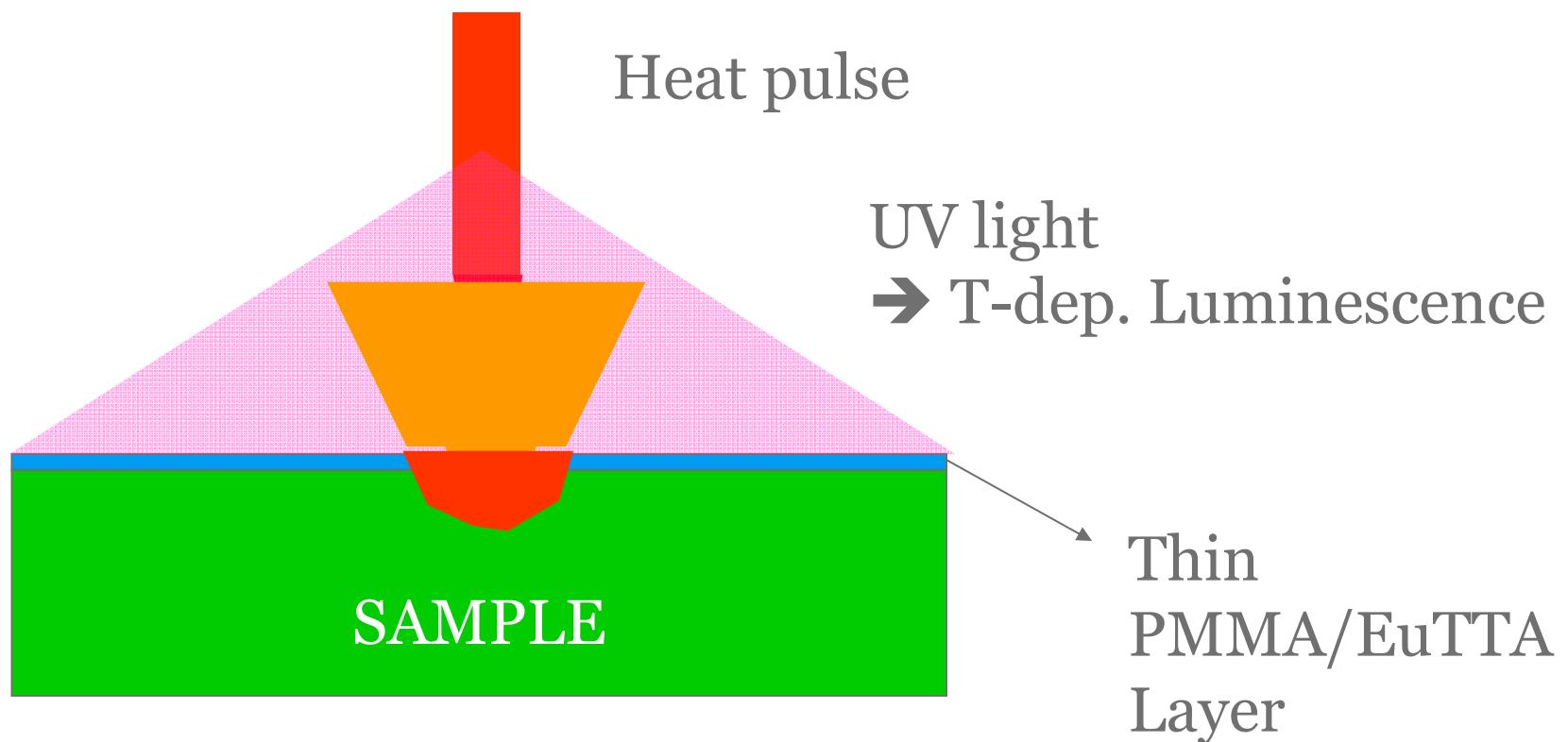
- › Heat source: Laser
- › Temperature detector: T-dependent luminescence

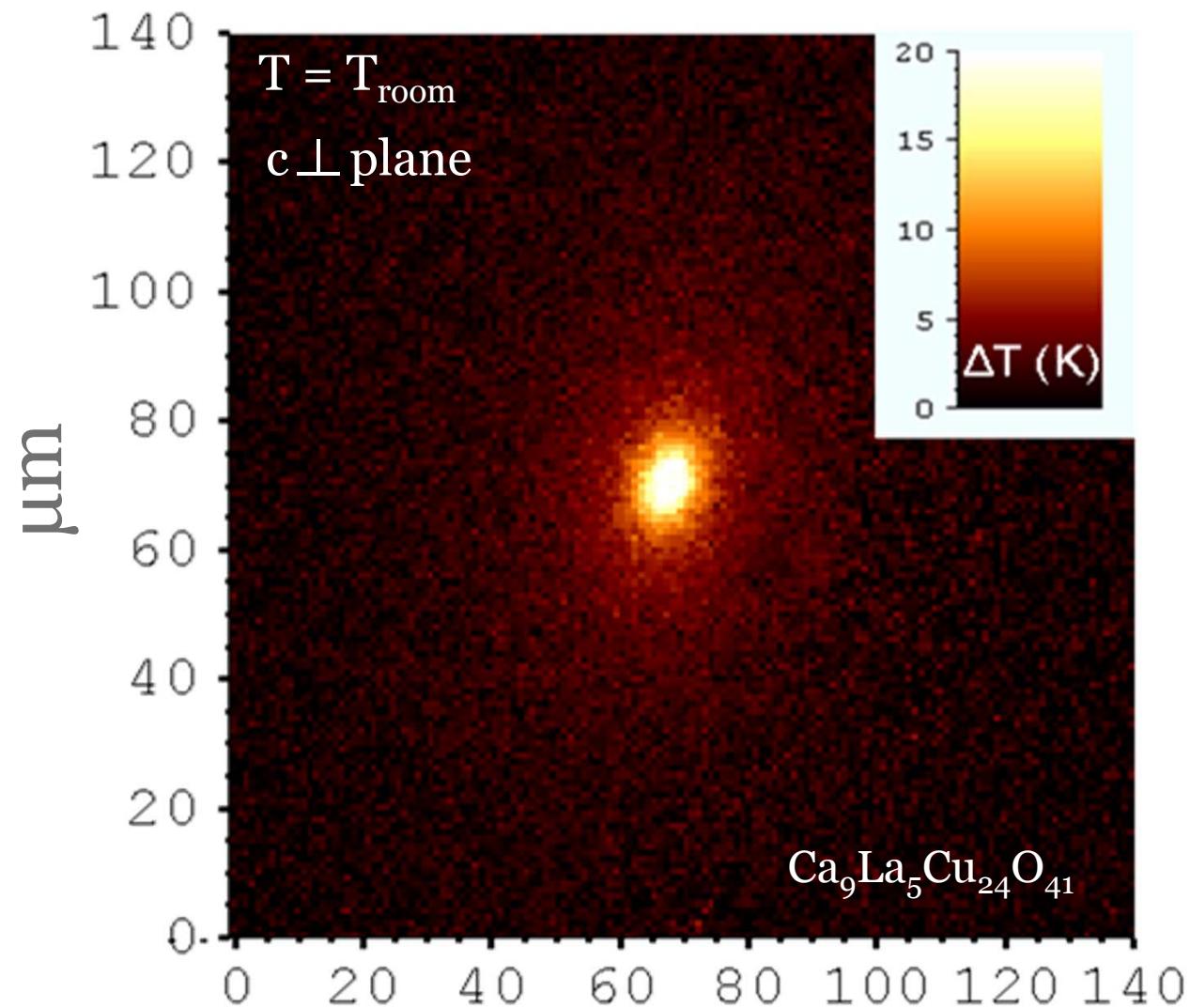


EuTTA: Eu-thienyltrifluoroacetonate



→ Spin coat EuTTA + dPMMA layer on crystals

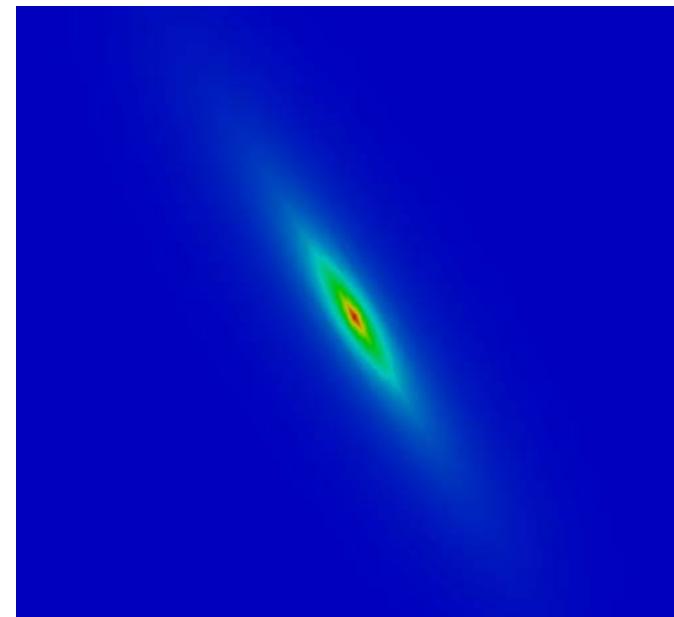
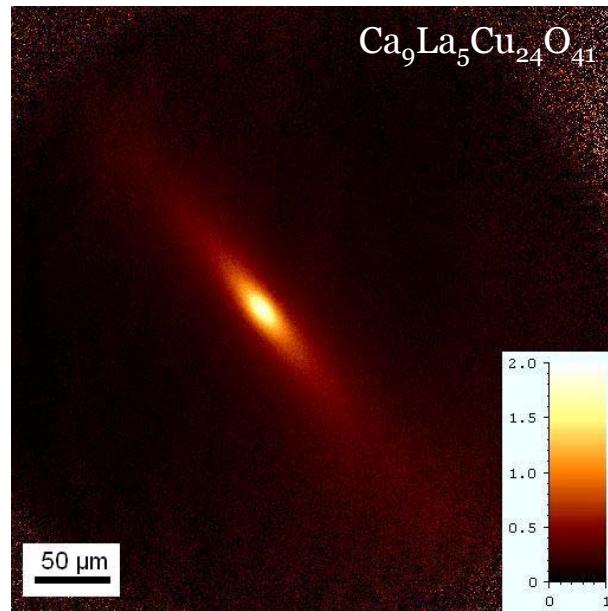




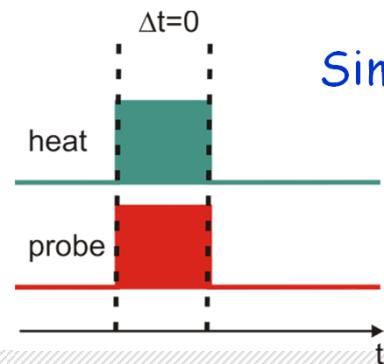


Static thermal imaging: ac-plane

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Experiment (RUG)



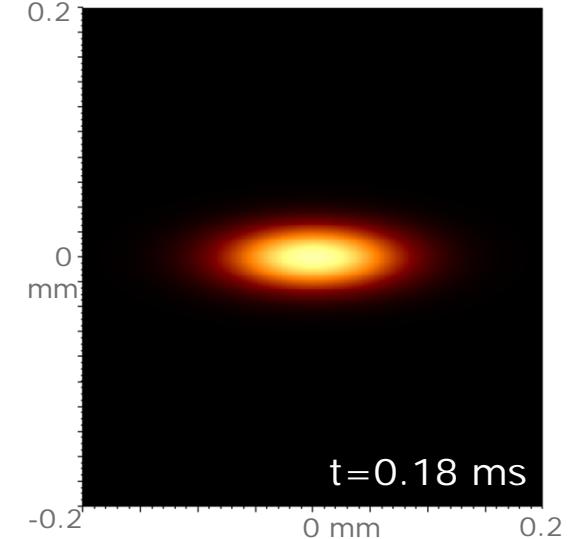
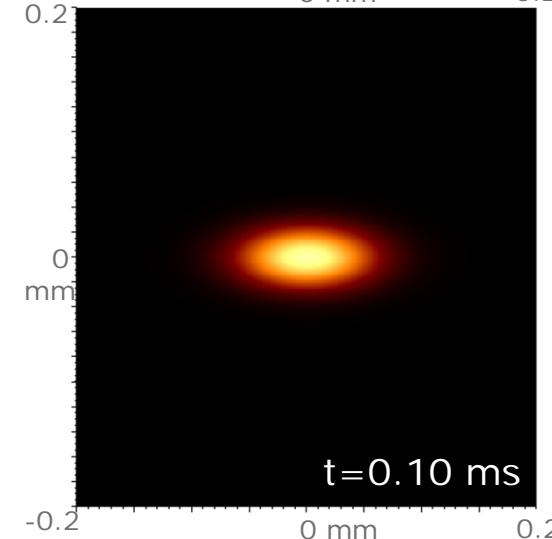
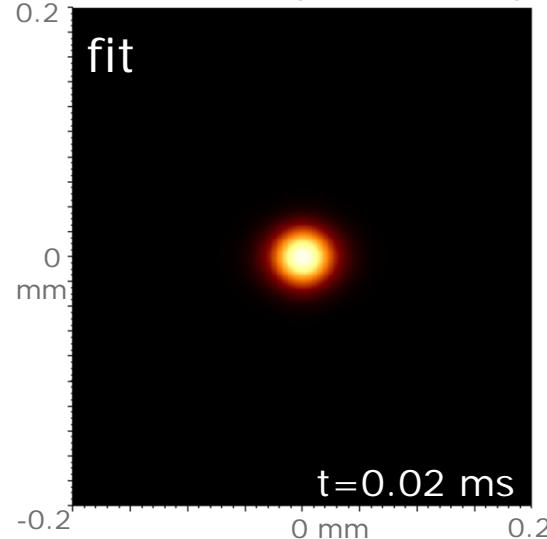
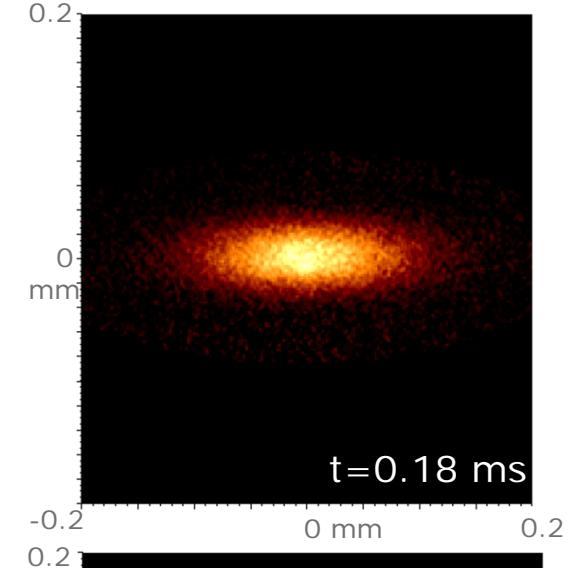
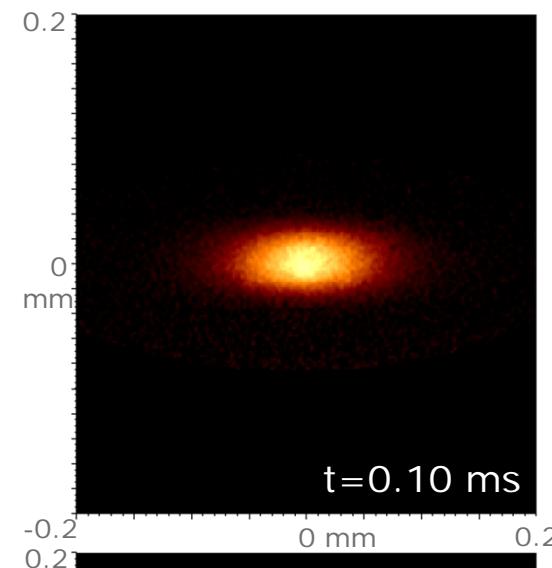
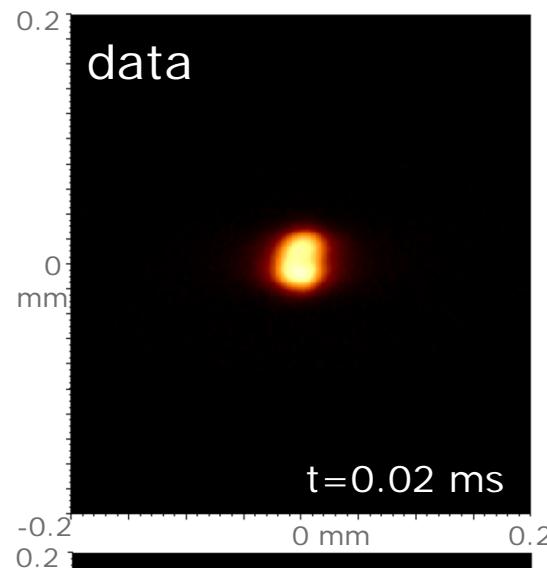
Simulation (ASCOMP)

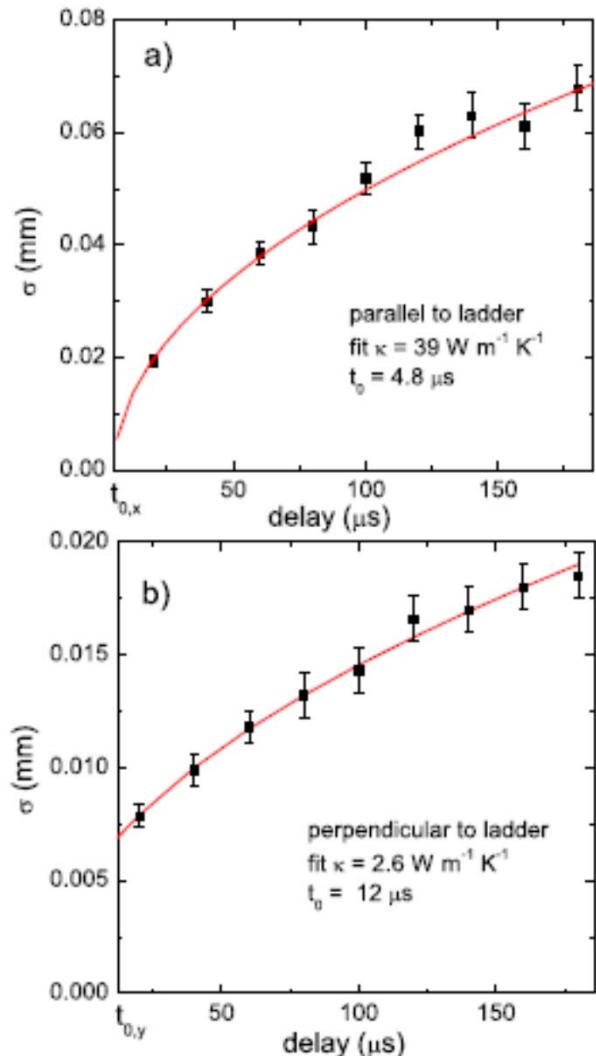
$$\frac{\partial T}{\partial t} = \nabla \left(\vec{D} \nabla \cdot T \right)$$



Heat transport in $\text{La}_5\text{Ca}_9\text{Cu}_{24}\text{O}_{41}$

| 32





Analysis using diffusion model

$$\text{Width: } \sigma_v = \sqrt{2D_v \cdot (t - t_{0,v})}$$
$$\kappa = \rho \cdot C_v \cdot D$$

Perpendicular to ladder
 $\kappa = 2.6 \text{ W/m}\cdot\text{K}$

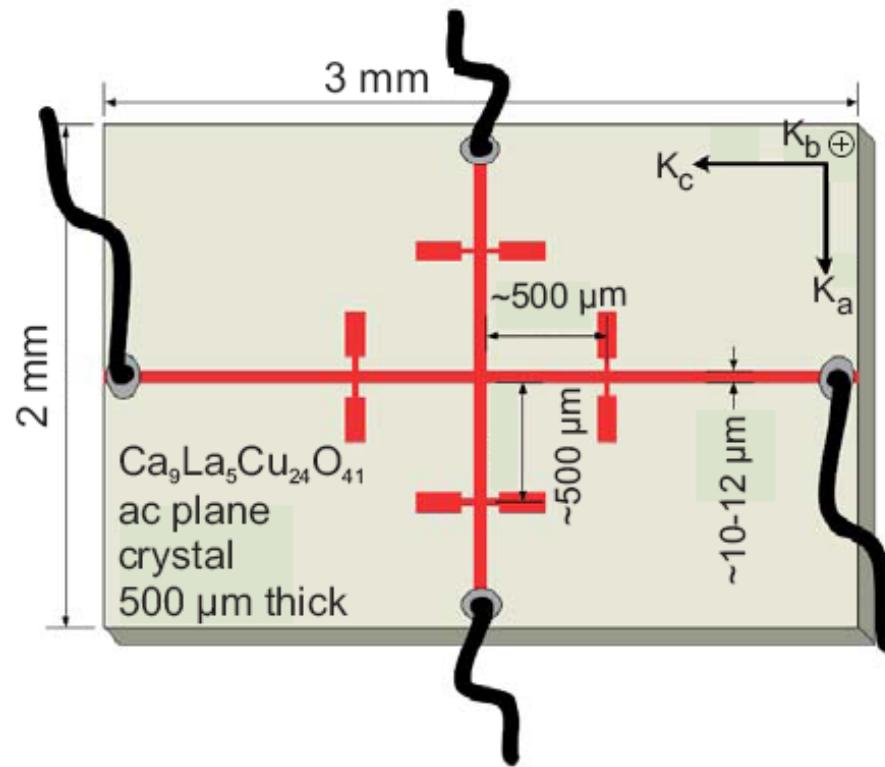
Good agreement with DC

Parallel to ladder
 $\kappa = 39 \text{ W/m}\cdot\text{K}$

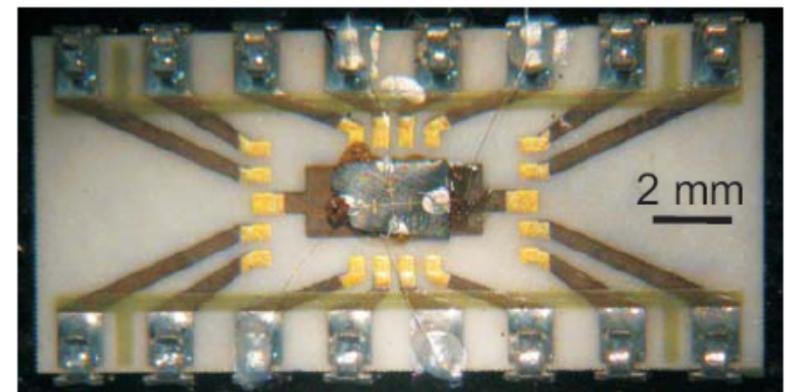
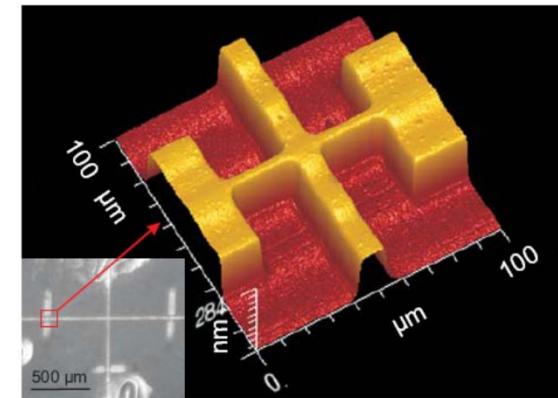
Bad agreement with DC
DC: $\kappa = 80 \text{ W/m}\cdot\text{K}$



Application in a device geometry



Hot strips along
a and c direction





Application in a device geometry

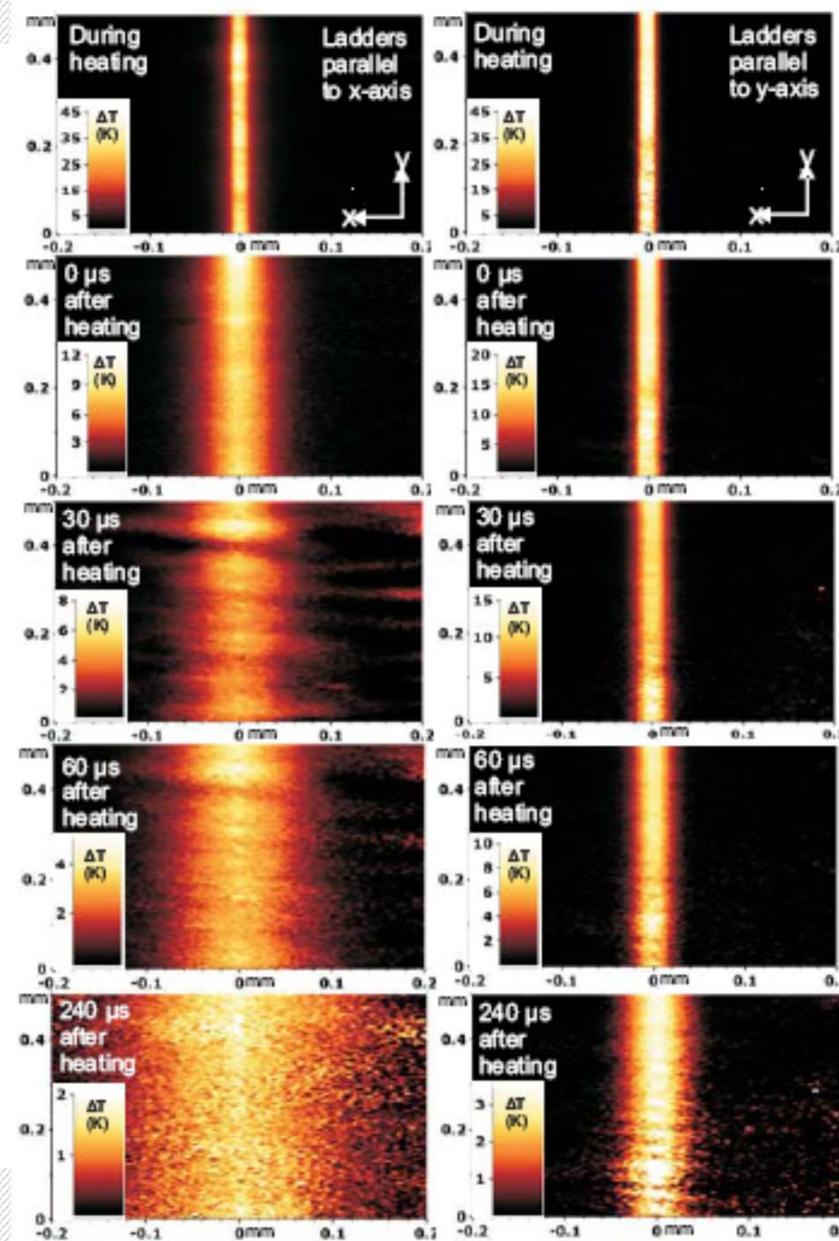
| 35

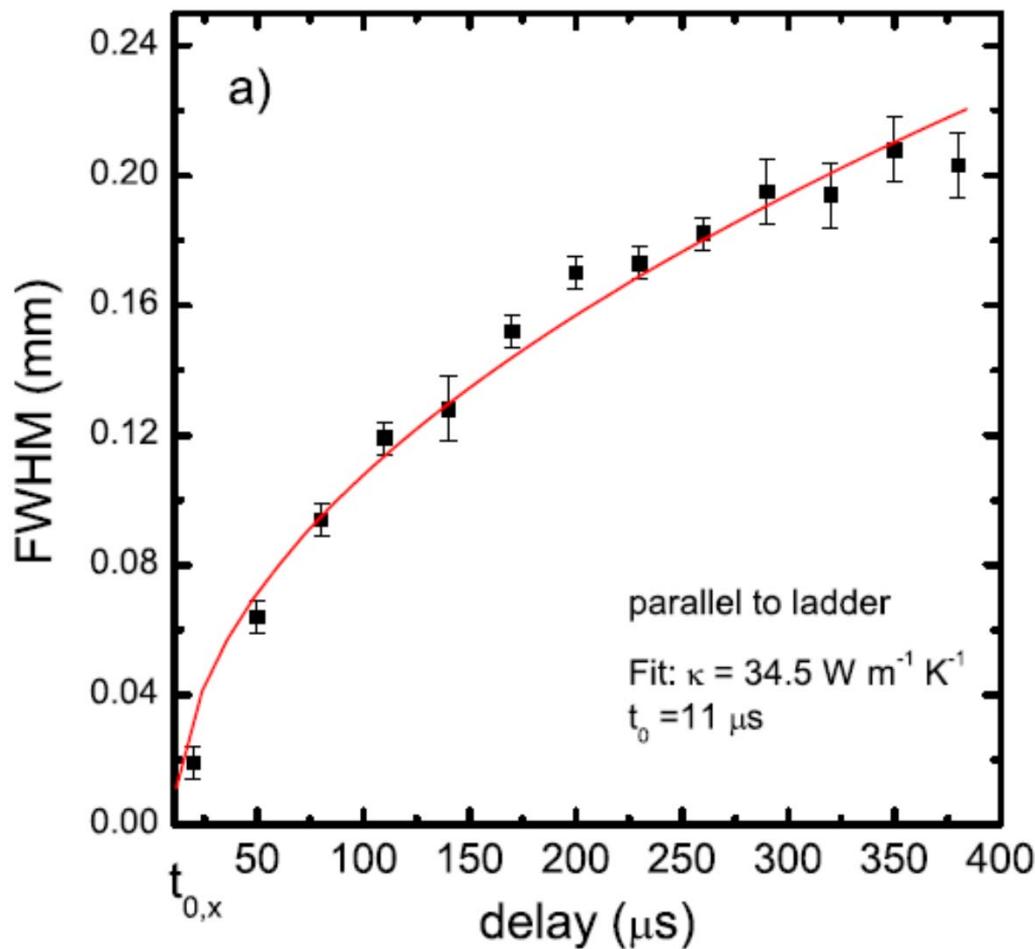
Strip perp. to ladder

Good
(magnon) conduction
along ladder

Strip along ladder

Bad
(phonon) conduction
perp to ladder





Heat conductivity

- Along a: 2.8 W/mK

- Along ladder (c-dir)

34.5 W/mK

Again: Discrepancy with DC !!!

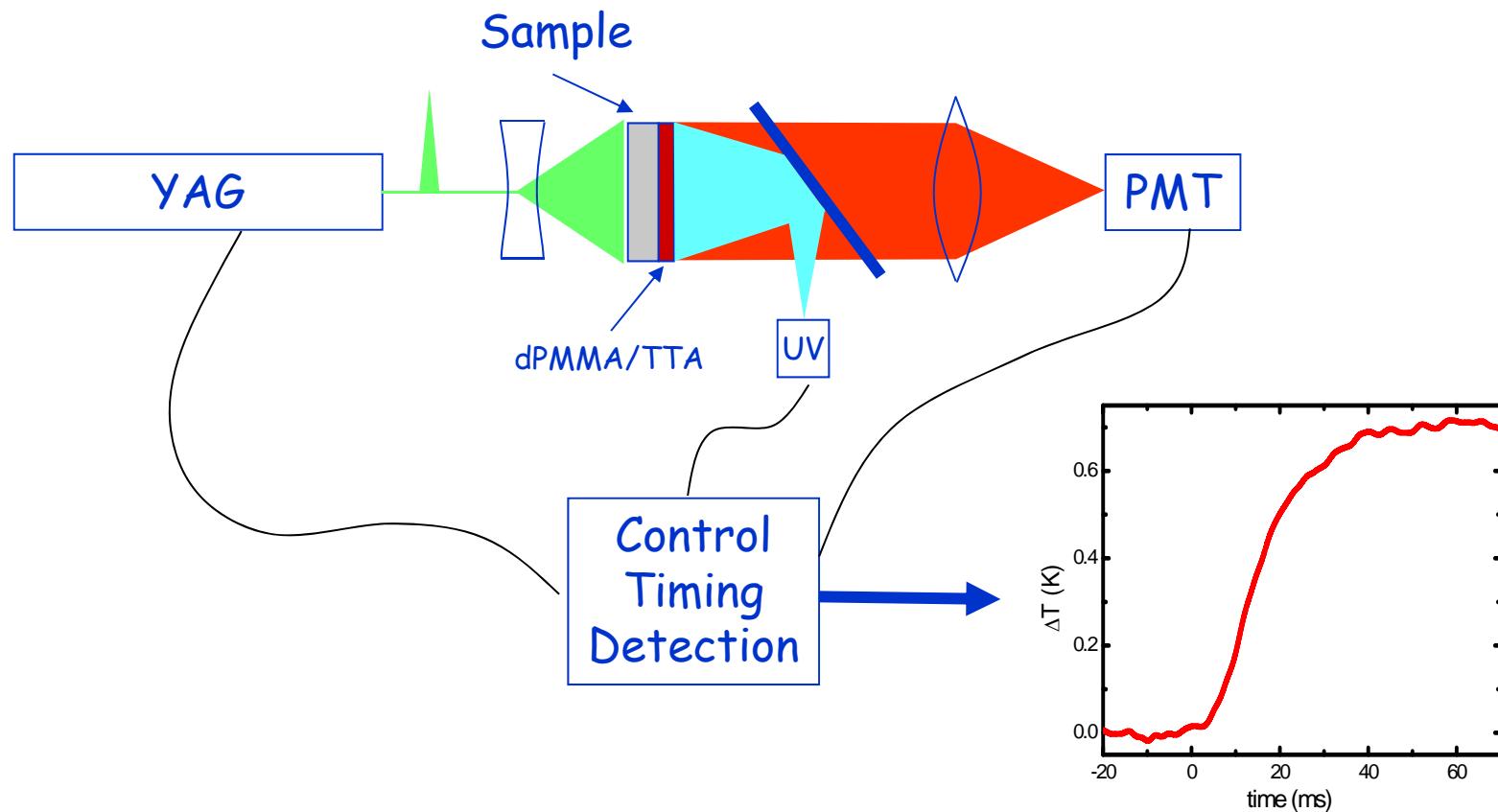


Discrepancy maybe due to surface effects?

Let us look at the bulk transport, again using optics



Heat diffusion set-up



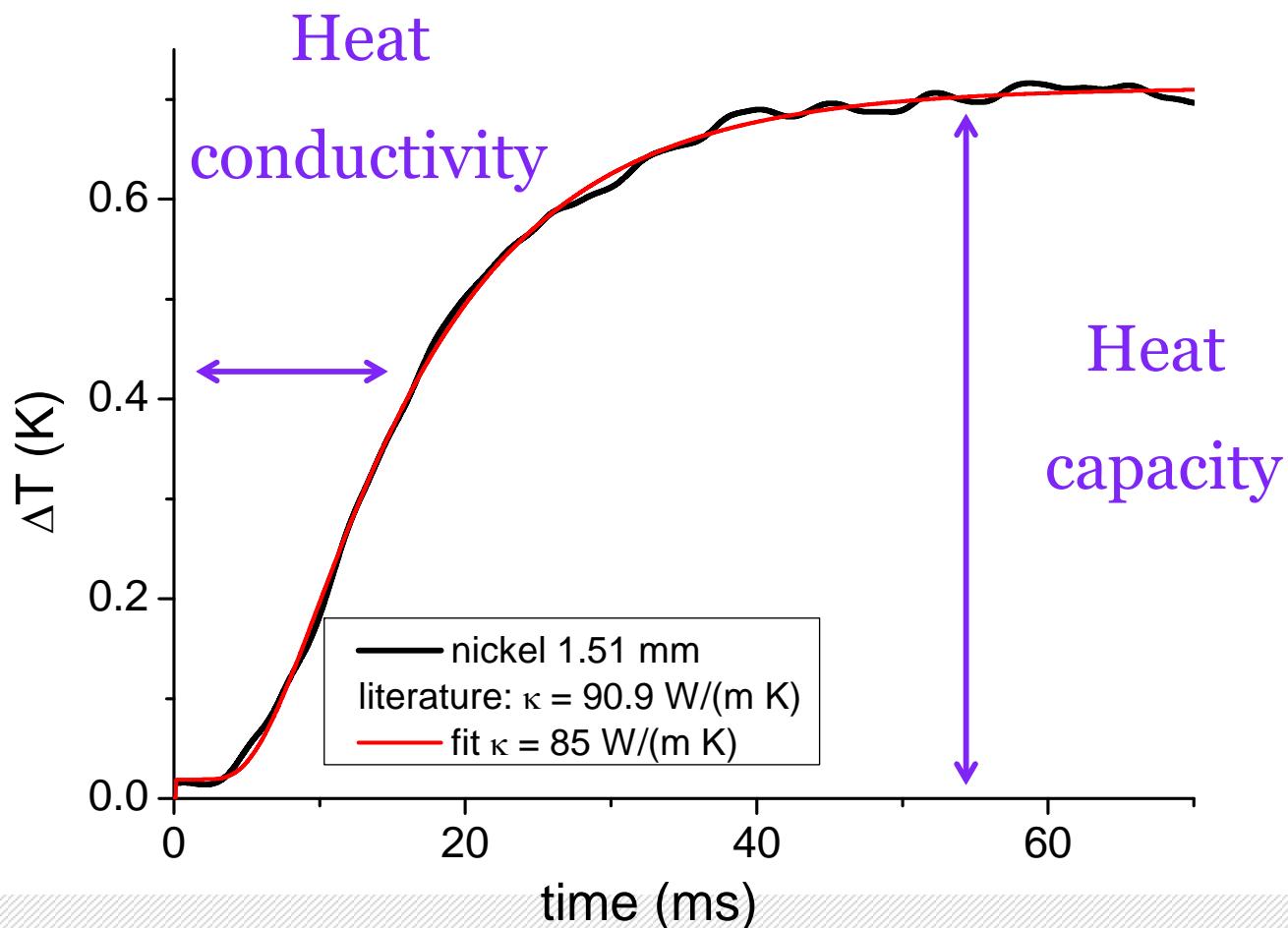


Test experiment on Ni single Xtal

| 39

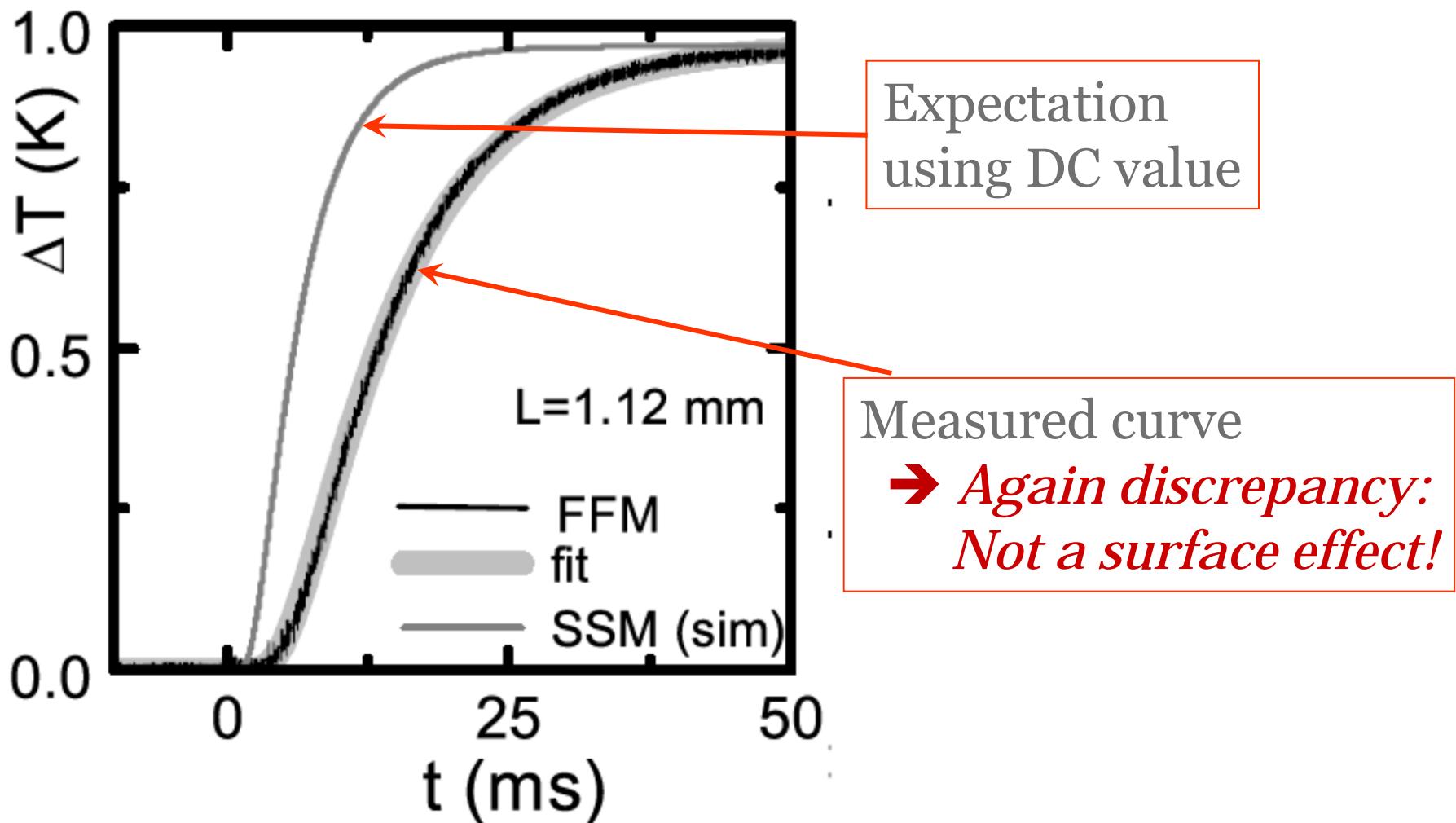
W.J. Parker et al, J. Appl. Phys. **32**, 1679 (1961).

$$\Delta T(t) = 1 + 2 \sum_{m=1}^{\infty} (-1)^m \exp(-m^2 \pi^2 D t / L^2)$$



$$\frac{\partial T}{\partial t} = D \frac{\partial^2 T}{\partial z^2}$$

$$D = \frac{\kappa}{\rho \cdot C_p}$$





We would like to probe magnon transport

One needs magnons for this

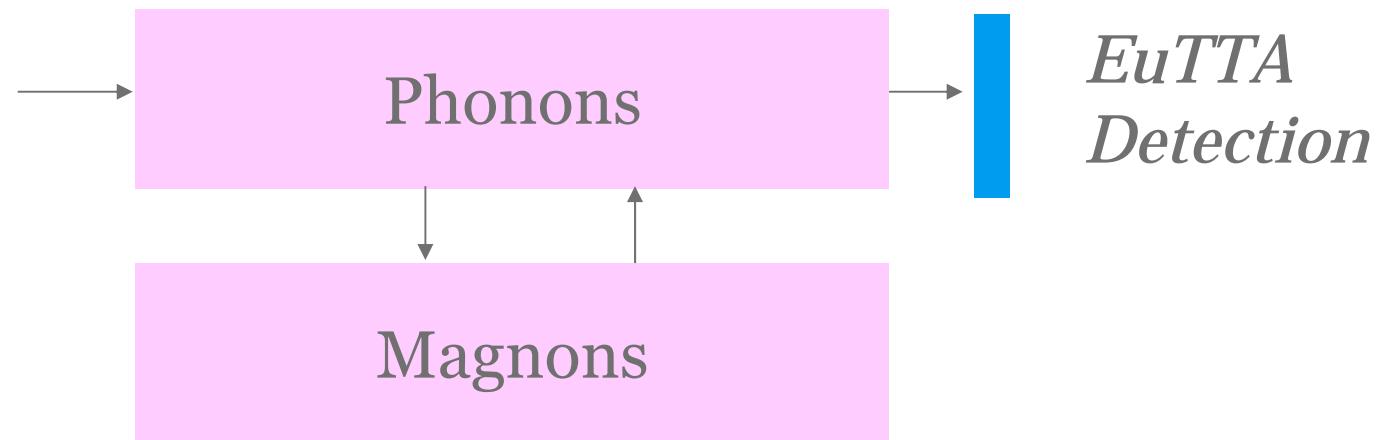
The experiments inject and detects phonons, not magnons !!!

In DC experiments, lots of time for equilibration

Optical experiment = Dynamic experiment !!!



Heat in (optical, Joule heating)



How fast is the conversion into magnons ???



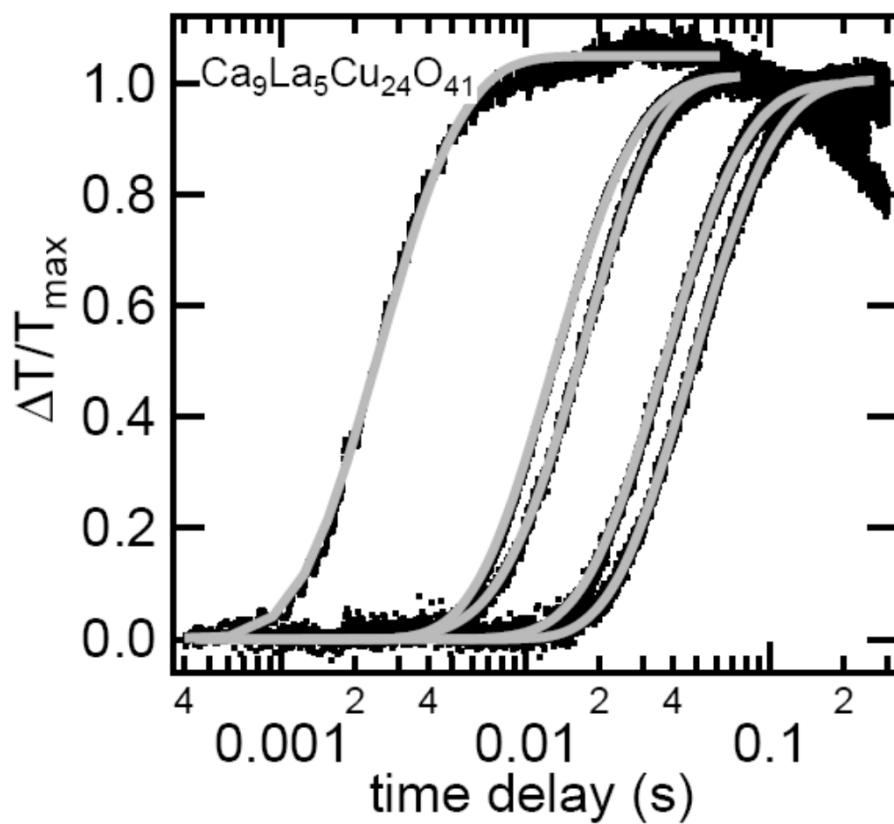
Coupled diffusion equations

$$C_L \partial_t T_L = \kappa_L \partial_x^2 T_L - g(T_L - T_M)$$

$$C_M \partial_t T_M = \kappa_M \partial_x^2 T_M + g(T_L - T_M)$$

$$g = \frac{C_L C_M}{C_L + C_M} \cdot \tau_{LM}$$

Solution $T(x=L, t)$ depends on thickness sample



Room T

Thickness l.t.r.

0.4 mm

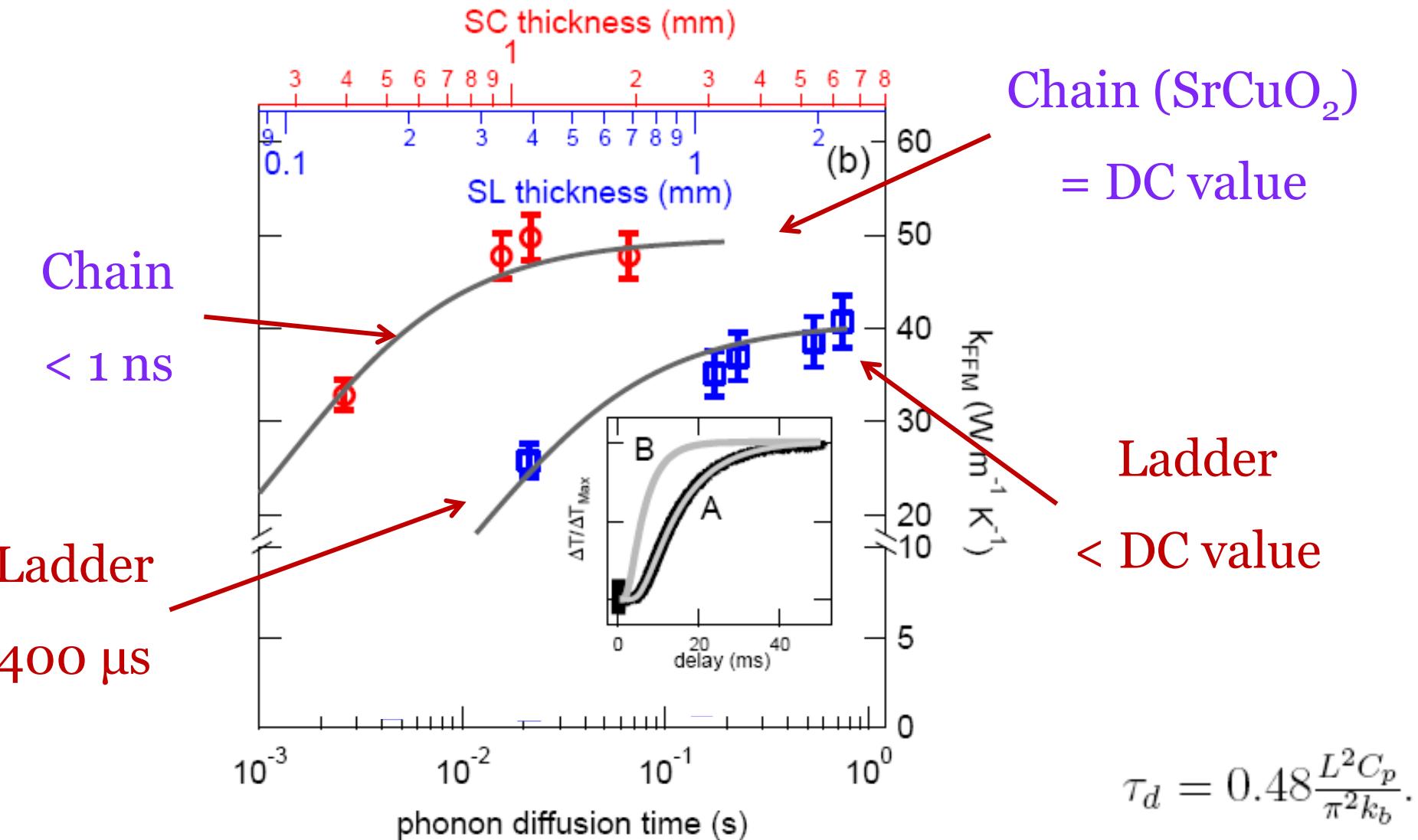
1.1

1.25

1.9

2.3

Fits to a 3D
two particle heat
transport model





- › Chain material: No large differences with DC xpt
- › Ladder material
 - Equilibration time: 300 μ s
 - ‘saturation’ value \sim half the DC value



PHYSICAL REVIEW B

VOLUME 15, NUMBER 3

1 FEBRUARY 1977

Effect of magnon-phonon thermal relaxation on heat transport by magnons*

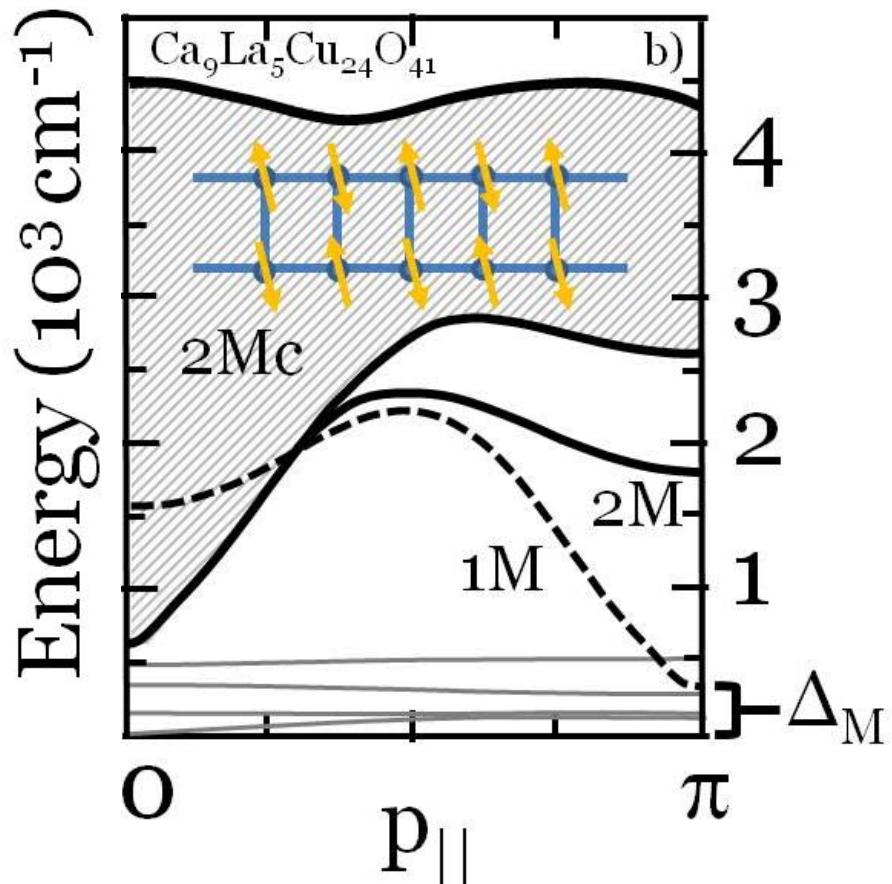
D. J. Sanders[†] and D. Walton

Department of Physics, McMaster University, Hamilton, Ontario, Canada

(Received 24 June 1976)

Alternatively, it should be noted³⁴ that $\tau_{m,p}$ varies as B_2^2 so that $\tau_{m,p}$ for MnF_2 should be roughly 100 times longer than that for YIG or $\sim 10^{-4}$ sec.

For MnF_2 phonon-magnon scattering time is estimated to be about 100 μs

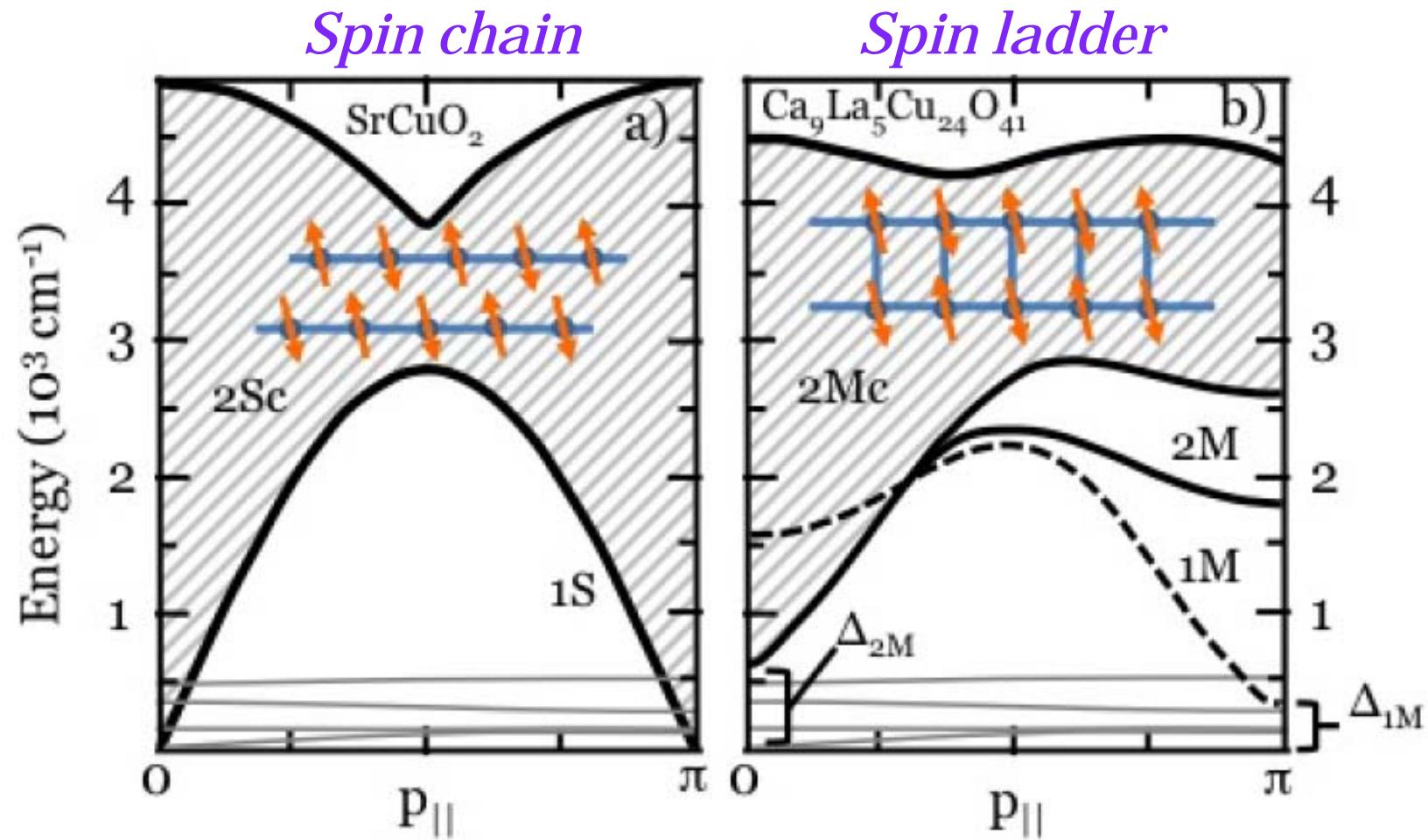


Heat mostly injected as acoustic phonons

Magnon gap $\Delta_M \sim 450 \text{ K}$

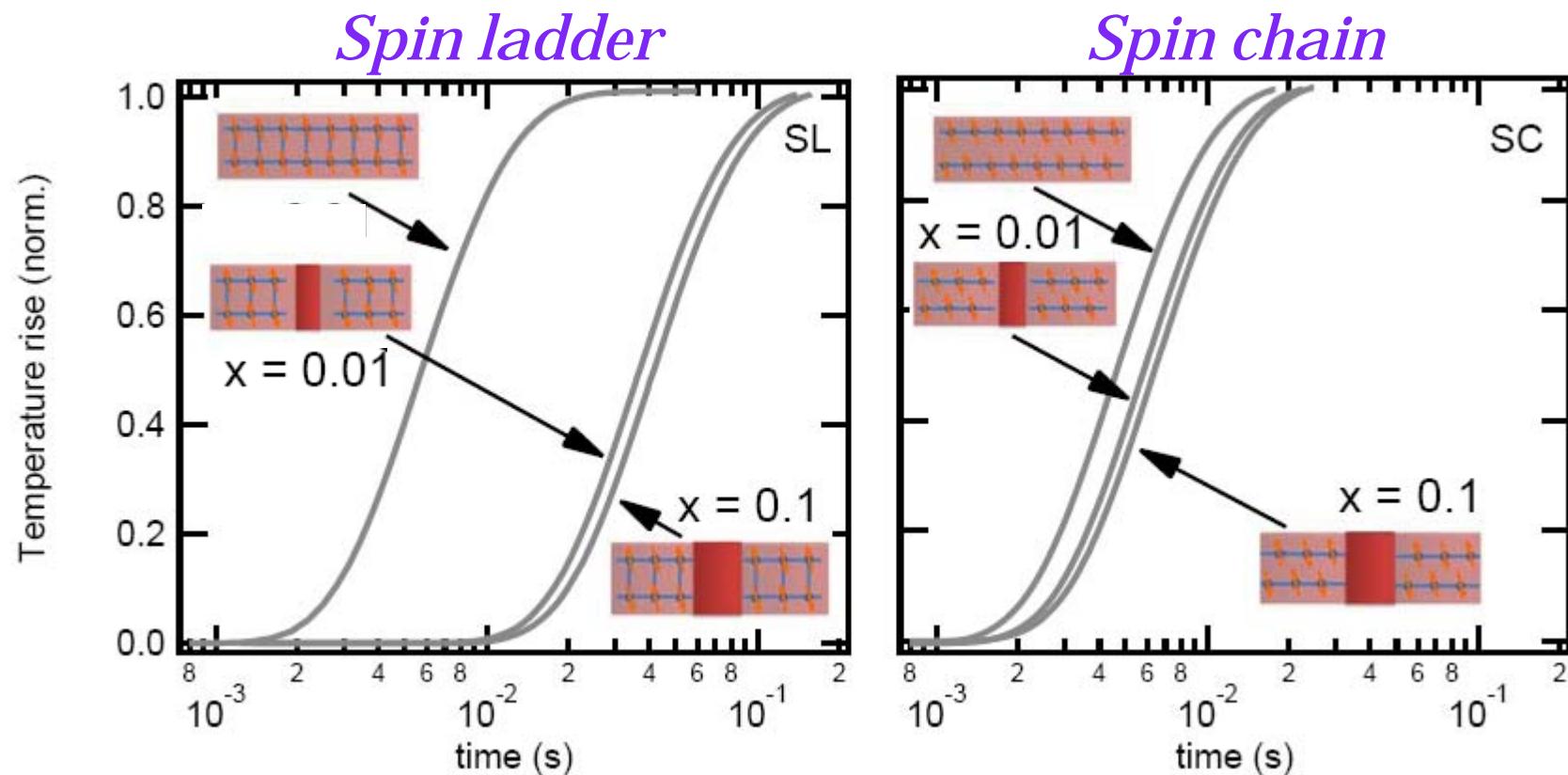
Phonon-magnon conversion
violates

- Energy conservation
- Spin conservation
- Momentum conservation





Defect (finite spin chains/ladders) limited Dynamic conductivity





Quantum magnets for heat transport

- › High thermal conductivity
- › uni-directional
- › electrically insulating

Diffusion and thermal imaging provide all-optical heat transport detection methods

- › straightforward with a fast response time
- › bonus: heat capacity
- › applicable to bulk and thin films and easy to integrate in pump-probe methods

Dynamic heat transport

- › Allows for a direct measurement of phonon-spin interaction
- › Coupling in ladder material very weak due to general conservation laws (E, k, S)



RUG

Marian Otter
Matteo Montagnese
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Thank you for your attention